

Observation of the decays $B^+ \rightarrow \Sigma_c(2455)^+ \bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0 \bar{\Xi}_c'^0$

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Li[✉], W. Z. Li[✉], Y. Li[✉], Y. B. Li[✉], Y. P. Liao[✉], J. Libby[✉], J. Lin[✉], S. Lin[✉], Z. Liptak[✉], V. Lisovskyi[✉], C. Liu[✉], M. H. Liu[✉], Q. Y. Liu[✉], Z. Q. Liu[✉], D. Liventsev[✉], S. Longo[✉], A. Lozar[✉], T. Lueck[✉], C. Lyu[✉], J. L. Ma[✉], Y. Ma[✉], M. Maggiora[✉], S. P. Maharana[✉], R. Maiti[✉], G. Mancinelli[✉], R. Manfredi[✉], E. Manoni[✉], M. Mantovano[✉], D. Marcantonio[✉], S. Marcello[✉], M. Marfoli[✉], C. Marinas[✉], C. Martellini[✉], A. Martens[✉], T. Martinov[✉], L. Massaccesi[✉], M. Masuda[✉], T. Matsuda[✉], D. Matvienko[✉], S. K. Maurya[✉], M. Maushart[✉], J. A. McKenna[✉], Z. Mediankin Gruberová[✉], R. Mehta[✉], F. Meier[✉], D. Meleshko[✉], M. Merola[✉], C. Miller[✉], M. Mirra[✉], K. Miyabayashi[✉], H. Miyake[✉], R. Mizuk[✉], S. Moneta[✉], A. L. Moreira de Carvalho[✉], H.-G. Moser[✉], N. Mudgal[✉], Th. Muller[✉], H. Murakami[✉], R. Mussa[✉], M. Nakao[✉], Y. Nakazawa[✉], Z. Natkaniec[✉], A. Natochii[✉], M. Nayak[✉], M. Neu[✉], M. Niiyama[✉], S. Nishida[✉], R. Nomaru[✉], S. Ogawa[✉], R. Okubo[✉], H. Ono[✉], G. Pakhlova[✉], S. Pardi[✉], J. Park[✉], K. Park[✉], S.-H. Park[✉], A. Passeri[✉], S. Patra[✉], T. K. Pedlar[✉], L. E. Piilonen[✉], P. L. M. Podesta-Lerma[✉], T. Podobnik[✉], L. Polat[✉], A. Prakash[✉], R. pramanik[✉], V. Prasad[✉], S. Prell[✉], E. Prencipe[✉], M. T. Prim[✉], H. Purwar[✉], P. Rados[✉], S. Raiz[✉], K. Ravindran[✉], J. U. Rehman[✉], M. Reif[✉], S. Reiter[✉], M. Remnev[✉], L. Reuter[✉], D. Ricalde Herrmann[✉], I. Ripp-Baudot[✉], G. Rizzo[✉], S. H. Robertson[✉], J. M. Roney[✉], A. Rostomyan[✉], N. Rout[✉], G. Russo[✉], S. Saha[✉], G. Sanchez[✉], D. A. Sanders[✉], S. Sandilya[✉], L. Santelj[✉], C. Santos[✉], V. Savinov[✉], B. Scavino[✉], J. Schmitz[✉], S. Schneider[✉], G. Schnell[✉], K. Schoenning[✉], C. Schwanda[✉], Y. Seino[✉], K. Senyo[✉], J. Serrano[✉], C. Sfienti[✉], W. Shan[✉], C. P. Shen[✉], X. D. Shi[✉], T. Shillington[✉], T. Shimasaki[✉], J.-G. Shiu[✉], D. Shtol[✉], A. Sibidanov[✉], F. Simon[✉], J. B. 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We report the first observation of the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$, with significances of 6.4σ and 5.3σ , respectively, including systematic uncertainties. This analysis is based on data samples containing 771.6×10^6 $\Upsilon(4S)$ decays collected with the Belle detector at the KEKB collider and 520.6×10^6 $\Upsilon(4S)$ decays collected with the Belle II detector at the SuperKEKB collider. The branching fractions are measured to be $\mathcal{B}(B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-) = (1.69 \pm 0.31 \pm 0.12_{-0.56}^{+1.39}) \times 10^{-3}$ and $\mathcal{B}(B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0) = (1.28 \pm 0.32 \pm 0.10_{-0.22}^{+0.26}) \times 10^{-3}$, where the first and second uncertainties are statistical and systematic, respectively, and the third arises from the uncertainties in the absolute branching fractions of $\bar{\Xi}_c'^-$ and $\bar{\Xi}_c'^0$ decays. This result represents the first observation of B -meson decays into a pair of charmed baryon-antibaryon states belonging to the same $SU(3)$ flavor sextet.

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B -meson decays provide a unique laboratory to explore the interplay between weak and strong interactions in the nonperturbative regime of quantum chromodynamics. Two-body baryonic B decays proceed at tree level through W -exchange, W -annihilation, and internal W -emission topologies [1,2]. The factorizable W -exchange and W -annihilation contributions are expected to be helicity suppressed [3]. However, in contrast to mesonic B decays, the internal W -emission contribution is not necessarily color suppressed and can therefore play a dominant role in baryonic two-body decays [3–5]. The large mass of the B meson allows for decays into charmed baryon-antibaryon pairs. Over the past two decades, several such modes have been observed [6–9]. The branching fractions of B -meson decays into charmed baryon-antibaryon pairs are significantly larger than those of two-body decays containing one charmed baryon and one charmless baryon, and both exceed those of fully charmless baryonic decays, reaching values as large as $\mathcal{O}(10^{-3})$ [10].

Recently, the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ were observed for the first time, by the Belle and Belle II experiments [11], with branching fractions measured to be of the order of 10^{-4} . These modes proceed through internal W -emission topology. The $\bar{\Xi}_c'$ and $\bar{\Xi}_c$ baryons share the same valence-quark content but in a heavy quark and light diquark model differ in the spin configuration of the light diquark. Consequently, the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ are expected to proceed through the same internal W -emission topology, as illustrated in Fig. 1. A comparison between the branching fractions of $B \rightarrow \Sigma_c(2455)\bar{\Xi}_c'$ and $B \rightarrow \Sigma_c(2455)\bar{\Xi}_c$ could provide insight into the underlying dynamics of B decays to charmed baryon-antibaryon pairs

belonging to different $SU(3)$ flavor multiplets. To date, there have been no observations of B decays into charmed baryon-antibaryon pairs in which both daughters belong to a flavor sextet.

In this paper, we report the first measurement of the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$. Most of the selection criteria follow those in Ref. [11]. In addition there are further requirements related to the additional photon from the $\bar{\Xi}_c'$ decay. The $\Sigma_c(2455)^{+,0}$ baryons are reconstructed via their decays into $\Lambda_c^+\pi^\pm$, with the Λ_c^+ reconstructed using the $\Lambda_c^+ \rightarrow pK^-\pi^+$ and $\Lambda_c^+ \rightarrow pK_S^0$ modes. The $\bar{\Xi}_c'^{0,-}$ baryons are reconstructed via their decays into $\gamma\bar{\Xi}_c'^{0,-}$, with the $\bar{\Xi}_c'$ reconstructed using the $\bar{\Xi}_c'^- \rightarrow \bar{\Xi}^+\pi^-\pi^-$ and $\bar{\Xi}_c'^- \rightarrow \bar{p}K^+\pi^-$ modes and the $\bar{\Xi}_c'^0$ reconstructed using the $\bar{\Xi}_c'^0 \rightarrow \bar{\Xi}^+\pi^-$ and $\bar{\Xi}_c'^0 \rightarrow \bar{\Lambda}K^+\pi^-$ modes. Charge-conjugate channels are implicitly included throughout this paper. This study is based on data samples containing $(771.6 \pm 10.6) \times 10^6$ and $(520.6 \pm 7.6) \times 10^6$ $\Upsilon(4S)$ events collected with the Belle [12] and Belle II [13] detectors, respectively, at the e^+e^- center-of-mass energy ($\sqrt{s}$) of 10.58 GeV, corresponding to integrated luminosities of 711 [14] and 492 fb^{-1} [15]. The analysis strategy is to fully reconstruct the signal decay chains and extract the signal yields from simultaneous fits to the distributions of the difference between the reconstructed B -meson energy and the beam energy at the $\Upsilon(4S)$ rest frame, both in the $\bar{\Xi}_c'$ signal and sideband regions.

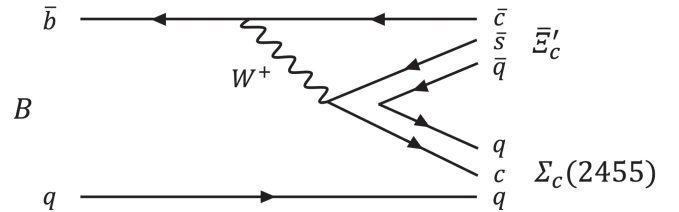


FIG. 1. Diagram illustrating the internal W -emission topology for the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$, corresponding to $q = u$ and $q = d$, respectively.

The Belle detector was a cylindrical large-solid-angle magnetic spectrometer located at the interaction point (IP) of the KEKB asymmetric-energy e^+e^- collider [16]. The detector consisted of a silicon vertex detector (SVD), a central drift chamber (CDC), an array of aerogel threshold Cherenkov counters (ACC), a barrel-like arrangement of time-of-flight (TOF) scintillation counters, and an electromagnetic calorimeter (ECL) composed of CsI(Tl) crystals located inside a superconducting solenoid coil that provided a 1.5 T axial magnetic field. An iron flux return located outside the coil was instrumented to detect K_L^0 mesons and to identify muons. A detailed description of the Belle detector can be found in Ref. [12].

The Belle II detector, located at the SuperKEKB asymmetric-energy e^+e^- collider [17], is an upgraded version of the Belle detector. Belle II includes a SVD, consisting of pixel sensors and double-sided strip detectors, and a CDC. The CDC is surrounded by two types of Cherenkov light detector systems: time-of-propagation (TOP) detector for the barrel region and an aerogel ring-imaging Cherenkov (ARICH) detector for the forward end-cap region. The Belle ECL is reused in Belle II along with the solenoid and the iron flux-return yoke. However, the ECL readout system has been upgraded. The solenoid flux return is instrumented with resistive-plate chambers and plastic scintillator modules to detect muons and K_L^0 mesons. A detailed description of the Belle II detector can be found in Ref. [13].

Simulated signal events are used to optimize the selection criteria, calculate the reconstruction efficiencies, and determine the fit models. The EVTGEN [18] and PYTHIA [19,20] software packages are used to generate $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$ events with final-state radiation simulated by the PHOTOS software package [21]. In the simulation, one B meson is not constrained to decay into a specific final state, while the other decays into a signal mode. The simulation includes beam-induced background data overlay [22]. Inclusive simulated samples of $e^+e^- \rightarrow q\bar{q}$ ($q = u, d, s, c$) and $\Upsilon(4S) \rightarrow B\bar{B}$ are used to optimize the selection criteria and identify the background sources [23], corresponding to 6 times the integrated luminosity of the Belle data and 4 times that of the Belle II data. The KKMC [24] and PYTHIA [19,20] software packages are used to simulate the $e^+e^- \rightarrow q\bar{q}$ processes. The detector responses are modeled by the software packages GEANT3 [25] for Belle and GEANT4 [26] for Belle II.

We analyze both Belle and Belle II data using the Belle II analysis software framework [27]. For the Belle data, the event reconstruction up to the particle level is performed with the Belle software, and the reconstructed data are then converted into a Belle II compatible format [28]. This allows the Belle and Belle II samples to be analyzed within a unified framework. To avoid experimental bias, the signal region was not examined until the

analysis procedure was finalized. All selection criteria are determined by iteratively optimizing the Punzi figure of merit for a target significance of 5 standard deviations in simulation [29]. We apply identical event-selection criteria in the Belle and Belle II analyses, unless otherwise stated.

We select tracks, except for those from the decays of K_S^0 , $\bar{\Lambda}$, and $\bar{\Xi}^+$, by requiring the impact parameters relative to the e^+e^- IP to be less than 2.0 cm perpendicular to the z -axis and less than 4.0 cm parallel to it. The z -axis is defined as the central solenoid axis with the positive direction in the e^- beam direction, common to both Belle and Belle II. The identification of charged tracks uses the likelihood ratio $\mathcal{R}(h|h') = \mathcal{L}(h)/[\mathcal{L}(h) + \mathcal{L}(h')]$, where $\mathcal{L}(h^{(i)})$ is the likelihood of the charged track being a hadron $h^{(i)} = p, K$, or π . This likelihood ratio is determined with a particle identification (PID) algorithm that combines information from different detector subsystems, including the specific ionization in the CDC, time-of-flight measurement in the TOF (TOP), and the response of the ACC (ARICH) in Belle (Belle II) [30,31]. Correction factors derived from control samples in data are applied to account for differences between data and simulation in tracking and PID efficiencies. Tracks with $\mathcal{R}(p|K) > 0.6$ and $\mathcal{R}(p|\pi) > 0.6$ are identified as proton candidates; charged kaon (pion) candidates must satisfy $\mathcal{R}(K|\pi) > 0.6$ (< 0.4). The efficiencies of these requirements range from 80% to 95%, with corresponding misidentification rates between 3% and 9%. The pion candidates used to reconstruct the K_S^0 , $\bar{\Lambda}$, and $\bar{\Xi}^+$ candidates are exempt from PID requirements as their kinematics provide sufficient discrimination.

The K_S^0 candidates are reconstructed from pairs of oppositely charged tracks treated as pions. An artificial neural network [32] and a boosted decision tree (BDT) classifier [33] are used in Belle and Belle II, respectively, to select the K_S^0 candidates. The $\bar{\Lambda}$ candidates are reconstructed from the decay $\bar{\Lambda} \rightarrow \bar{p}\pi^+$ and selected using momentum-dependent requirements on variables describing the displaced decay topology in Belle [34] and using a BDT classifier in Belle II. These discriminators primarily rely on the kinematic properties of the K_S^0 or $\bar{\Lambda}$ and their decay products. The invariant masses of the K_S^0 and $\bar{\Lambda}$ candidates are required to be within ± 9.0 and ± 5.5 MeV/ c^2 of their known masses [10], respectively, corresponding to approximately 2.5 times the mass resolution (σ). The selected $\bar{\Lambda}$ candidate is then combined with a π^+ candidate to form a $\bar{\Xi}^+$ candidate. The invariant mass of $\bar{\Xi}^+$ candidates is required to be within ± 6.5 MeV/ c^2 of its known mass [10] (approximately 2.5σ). A vertex fit is performed for each K_S^0 , $\bar{\Lambda}$, and $\bar{\Xi}^+$ candidate, and the candidate mass is constrained to the corresponding known value [10].

We define a photon candidate as an energy deposit in the ECL not associated with the extrapolated trajectory of any

charged track. To suppress backgrounds from neutral hadrons, we require $E(3 \times 3)/E(5 \times 5) \geq 0.6$, where $E(n \times n)$ denotes the energy deposited in an $n \times n$ array of crystals centered on the highest-energy crystal (for Belle II, the four corner crystals of the 5×5 array are excluded). The photon energy is required to exceed 50 MeV in the laboratory frame. At Belle II, BDT classifiers are used to suppress energy deposits from secondary hadronic interactions unrelated to the relevant collision and overlapping beam-induced photons [35]. The combined signal efficiencies (background rejection rates) of the two BDTs are approximately 89% (47%) and 89% (44%) for the charged and neutral channels, respectively.

The invariant masses of the reconstructed Λ_c^+ , Ξ_c^0 , and $\Xi_c^{\prime 0}$ candidates are required to lie within ± 15 , ± 18 , and ± 18 MeV/ c^2 of their known values [10], respectively, corresponding in each case to approximately 2.5σ . Each Λ_c^+ and $\Xi_c^{\prime 0}$ candidate is fitted to a common vertex, and the candidate mass is constrained to the corresponding known value [10]. The $\Sigma_c(2455)^{++0}$ baryons are reconstructed via their decays into $\Lambda_c^+\pi^\pm$, with the signal region defined as $2.446 < M(\Lambda_c^+\pi^\pm) < 2.464$ GeV/ c^2 . After applying this requirement, more than 95% of the signal events are retained. The sideband regions of $M(\Lambda_c^+\pi^\pm)$ are defined as $2.436 < M(\Lambda_c^+\pi^\pm) < 2.444$ and $2.466 < M(\Lambda_c^+\pi^\pm) < 2.494$ GeV/ c^2 . The $\Xi_c^{\prime -0}$ baryons are reconstructed via their decays into $\gamma\Xi_c^{\prime -0}$, with the reconstructed masses required to be within ± 20 MeV/ c^2 of the known $\Xi_c^{\prime -0}$ masses [10], corresponding in each case to approximately 2.5σ . Vertex fits are applied to the $\Sigma_c(2455)^{++0}$ and $\Xi_c^{\prime -0}$ candidates, with their masses constrained to the known values [10].

The B^+ and B^0 candidates are reconstructed from the $\Sigma_c(2455)^{++}\Xi_c^{\prime -}$ and $\Sigma_c(2455)^0\Xi_c^{\prime 0}$ combinations, respectively. This implies four distinct reconstruction modes for each signal channel. A vertex fit is applied to the B candidates with a requirement of $\chi^2/\text{ndf} < 20$, where ndf is the number of degrees of freedom. The beam-constrained mass M_{bc} and the energy difference ΔE are defined for B candidates as $M_{\text{bc}} = \sqrt{E_{\text{beam}}^2 - (\sum_i \vec{p}_i)^2}$ and $\Delta E = \sum_i E_i - E_{\text{beam}}$, where $E_{\text{beam}} = \sqrt{s}/2$ is the beam energy and $(E_i, \vec{p}_i)$ is the four-momentum of the i th daughter of the B meson, all calculated in the e^+e^- c.m. system. Candidates satisfying $M_{\text{bc}} > 5.27$ GeV/ c^2 (approximately 2.5σ) and $|\Delta E| < 0.06$ GeV (approximately 12σ) are retained for further analysis. The scatter plots of $M(\Lambda_c^+\pi^\pm)$ versus $M(\gamma\Xi_c^{\prime -0})$ for reconstructed $B^+ \rightarrow \Sigma_c(2455)^{++}\Xi_c^{\prime -}$ and $B^0 \rightarrow \Sigma_c(2455)^0\Xi_c^{\prime 0}$ candidates in the signal and sideband regions of ΔE from the combined Belle and Belle II data samples are shown in Appendix A, along with their one-dimensional projections onto $M(\Lambda_c^+\pi^\pm)$ and $M(\gamma\Xi_c^{\prime -0})$ in the signal region of ΔE . Because the photon from the $\Xi_c^{\prime -}$ decay has very low energy,

many random low-energy photon candidates, such as residual neutral clusters in the electromagnetic calorimeter or photons from other particles in the event, can form incorrect combinations with the $\Xi_c^{\prime -}$ candidates. This results in a large number of events with multiple candidates. After applying all the selection criteria, the fractions of events in signal simulation with multiple candidates are 60.5%–63.8% (54.3%–58.9%) for the charged channel and 61.0%–64.0% (53.8%–58.4%) for the neutral channel in Belle (Belle II), depending on the reconstruction mode. The corresponding average candidate multiplicities are 2.12–2.33 (1.93–2.13) for the charged channel and 2.15–2.29 (1.93–2.03) for the neutral channel in Belle (Belle II). For events with multiple candidates, a best-candidate selection (BCS) is applied, retaining the candidate with the smallest total χ^2 from the vertex and mass-constrained fits to the charmed baryons and the B meson. It was checked that the BCS procedure does not introduce peaking structures in the ΔE distribution or distort its shape, and yields signal efficiencies of 73.0%–76.2% (72.5%–75.1%) for the charged channel and 73.5%–74.0% (72.0%–75.0%) for the neutral channel in Belle (Belle II).

It is estimated from simulation that after the BCS around 50% of the selected signal candidates are incorrectly reconstructed, and these are referred to as self-cross-feed (SCF) events. These SCF events mainly originate from the reconstruction of the $\Xi_c^{\prime -0} \rightarrow \gamma\Xi_c^{\prime -0}$ decays, in which the genuine low-energy photon has been replaced by a random low-energy photon candidate. The SCF component exhibits a broad ΔE distribution that is clearly distinguishable from the correctly reconstructed signal but overlaps significantly with the combinatorial background. The SCF events are not treated as signal when determining the branching fractions. To have better control of SCF, we define the sideband regions of $M(\gamma\Xi_c^{\prime -0})$ as $2.508 < M(\gamma\Xi_c^{\prime -0}) < 2.548$ and $2.608 < M(\gamma\Xi_c^{\prime -0}) < 2.658$ GeV/ c^2 . The centers of the sideband regions are about 5σ away from the center of the signal region. The sideband regions contain a significant SCF contribution, with a shape in the ΔE distribution similar to that observed in the signal region. To better constrain the SCF component in the signal extraction, we perform a simultaneous fit to the ΔE distributions in the signal and sideband regions of $M(\gamma\Xi_c^{\prime -0})$. Backgrounds are studied using inclusive simulated samples [23] and data from the sidebands of $M(\Lambda_c^+)$, $M(\Lambda_c^+\pi^\pm)$, $M(\Xi_c^{\prime 0})$, and M_{bc} distributions. No peaking background is observed in the ΔE distributions of events from these sidebands.

To extract the signal yield of $B^+ \rightarrow \Sigma_c(2455)^{++}\Xi_c^{\prime -}$ or $B^0 \rightarrow \Sigma_c(2455)^0\Xi_c^{\prime 0}$, we perform simultaneous unbinned extended maximum-likelihood fits to the ΔE distributions of four samples: events from the signal and sideband regions of $M(\gamma\Xi_c^{\prime -0})$ in both Belle and Belle II data. The likelihood functions used to model the ΔE distributions in

the signal and sideband regions of $M(\gamma\bar{\Xi}_c^{\prime-0})$ are parametrized as

$$\mathcal{L} = \frac{e^{-(N_{\text{sig}}+N_{\text{scf}}+N_{\text{bkg}})}}{N!} \prod_i^N [N_{\text{sig}}\mathcal{F}_{\text{sig}}(\Delta E_i) + N_{\text{scf}}\mathcal{F}_{\text{scf}}(\Delta E_i) + N_{\text{bkg}}\mathcal{F}_{\text{bkg}}(\Delta E_i)]$$

and

$$\mathcal{L}^{\text{sb}} = \frac{e^{-(N_{\text{sig}}^{\text{sb}}+N_{\text{scf}}^{\text{sb}}+N_{\text{bkg}}^{\text{sb}})}}{N^{\text{sb}}!} \prod_i^{N^{\text{sb}}} [N_{\text{sig}}^{\text{sb}}\mathcal{F}_{\text{sig}}^{\text{sb}}(\Delta E_i) + N_{\text{scf}}^{\text{sb}}\mathcal{F}_{\text{scf}}^{\text{sb}}(\Delta E_i) + N_{\text{bkg}}^{\text{sb}}\mathcal{F}_{\text{bkg}}^{\text{sb}}(\Delta E_i)],$$

respectively. Here, $\mathcal{F}^{(\text{sb})}(\Delta E_i)$ denotes the probability density function (PDF) evaluated at the measured ΔE value of the i th event, and $N^{(\text{sb})}$ denotes the number of observed events. The signal PDF, $\mathcal{F}_{\text{sig}}$, is parametrized by a double-Gaussian function with the same mean values. All parameters of the double-Gaussian function are fixed to the values obtained from simulation. The SCF PDFs, $\mathcal{F}_{\text{scf}}^{(\text{sb})}$, are modeled using the nonparametric kernel density estimation [36] derived from simulation. The combinatorial

background PDFs, $\mathcal{F}_{\text{bkg}}^{(\text{sb})}$, are described by first-order polynomials with free parameters. The polynomial parameters are allowed to vary independently between the signal and sideband regions of $M(\gamma\bar{\Xi}_c^{\prime-0})$ and between the Belle and Belle II data samples. The quantities $N_{\text{sig}}^{(\text{sb})}$, $N_{\text{scf}}^{(\text{sb})}$, and $N_{\text{bkg}}^{(\text{sb})}$ represent the signal, SCF, and combinatorial background yields, respectively. The signal (SCF) yields satisfy $N_{\text{sig}}^{\text{sb}} = f_{\text{sig}} N_{\text{sig}}$ ($N_{\text{scf}}^{\text{sb}} = f_{\text{scf}} N_{\text{scf}}$), where $f_{\text{sig}} (f_{\text{scf}})$ is defined as the ratio of the signal (SCF) yield in the sideband regions to that in the signal region of $M(\gamma\bar{\Xi}_c^{\prime-0})$, as estimated from simulation. The values of f_{sig} and f_{scf} are fixed by simulation to 0.02 (0.05) and 1.9 (1.8) for the charged channel and 0.01 (0.04) and 1.8 (1.6) for the neutral channel in Belle (Belle II), respectively. The event yields are determined from the fit, with the signal yields in the Belle and Belle II datasets constrained to follow the expected ratio for a common branching fraction.

Figure 2 shows the ΔE distributions of candidates reconstructed in the signal and sideband regions of $M(\gamma\bar{\Xi}_c^{\prime-0})$ in the combined Belle and Belle II data, with fit results overlaid. The fitted signal yields for the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime-}$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime0}$ are 62.3 ± 11.3 and 30.7 ± 7.6 , respectively, with statistical

$$\int \mathcal{L} dt = 711 \text{ fb}^{-1} \text{ (Belle)} + 492 \text{ fb}^{-1} \text{ (Belle II)}$$

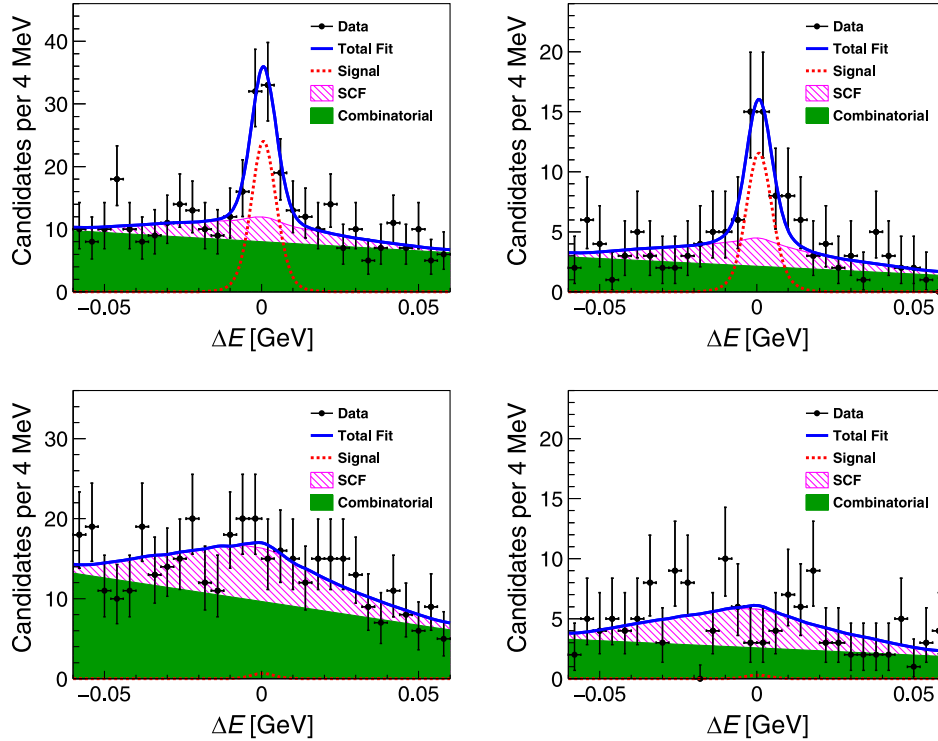


FIG. 2. Distributions of ΔE for the reconstructed (left) $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime-}$ and (right) $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime0}$ candidates in the (top) signal and (bottom) sideband regions of $M(\gamma\bar{\Xi}_c^{\prime-0})$ in the combined Belle and Belle II data. Points with error bars denote the data. The blue solid and red dashed curves represent the total fit results and fitted signal components, respectively, while the magenta hatched and green filled histograms represent the fitted SCF and combinatorial background components.

TABLE I. Summary of the fitted signal yields, the values of $\sum_i(\epsilon_i^{\text{B},\text{B}2}\mathcal{B}_i)$, and the measured branching fractions of $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ decays. We list only the statistical uncertainties of the signal yields. For the branching fractions, the first and second uncertainties are statistical and systematic, respectively, while the third originates from the absolute branching fractions of the $\bar{\Xi}_c'^0$ decays [10].

	N_{sig}	$\sum_i(\epsilon_i^{\text{B}}\mathcal{B}_i)(10^{-5})$	$\sum_i(\epsilon_i^{\text{B}2}\mathcal{B}_i)(10^{-5})$	$\mathcal{B}(10^{-3})$
B^+	62.3 ± 11.3	2.43	3.32	$1.69 \pm 0.31 \pm 0.12^{+1.39}_{-0.56}$
B^0	30.7 ± 7.6	1.53	2.49	$1.28 \pm 0.32 \pm 0.10^{+0.26}_{-0.22}$

significances of 8.6σ and 6.9σ . The statistical significances are calculated using $-2\ln(\mathcal{L}_0/\mathcal{L}_{\text{max}})$, accounting for the difference in the number of degrees of freedom ($\Delta\text{ndf} = 3$), where $\mathcal{L}_0$ and $\mathcal{L}_{\text{max}}$ are the likelihood values obtained from fits without and with the signal and SCF components, respectively.

To evaluate the signal significances including systematic uncertainties, we repeat the significance calculation using the same alternative fits as those used to evaluate the systematic uncertainties associated with the fit models, with the relevant variations considered in combination. The most conservative values, corresponding to the smallest significances obtained among all tested cases, are 6.4σ for the $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ channel and 5.3σ for the $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ channel. These values are taken as the final signal significances after incorporating systematic effects.

The branching fractions of $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ decays are calculated using

$$\mathcal{B} = \frac{N_{\text{sig}}}{2f_{+-/00}[N_{\Upsilon(4S)}^{\text{B}}\sum_i(\epsilon_i^{\text{B}}\mathcal{B}_i) + N_{\Upsilon(4S)}^{\text{B}2}\sum_i(\epsilon_i^{\text{B}2}\mathcal{B}_i)]}.$$

Here, N_{sig} represents the number of fitted $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-$ or $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0$ signal events in the combined Belle and Belle II datasets; $N_{\Upsilon(4S)}^{\text{B},\text{B}2}$ denotes the total number of $\Upsilon(4S)$ events in the Belle or Belle II datasets; $f_{+-/00}$ refers to the fraction of charged ($f_{+-} = 0.5113^{+0.0073}_{-0.0108}$) or neutral ($f_{00} = 0.4861^{+0.0074}_{-0.0080}$) $B\bar{B}$ pairs [37]; the term $\sum_i(\epsilon_i^{\text{B},\text{B}2}\mathcal{B}_i)$ represents the sum of the reconstruction efficiencies $\epsilon_i^{\text{B},\text{B}2}$ (for Belle or Belle II) times the corresponding secondary branching fractions $\mathcal{B}_i$ over all the reconstruction modes. The values of $\epsilon_i^{\text{B},\text{B}2}$ and $\epsilon_i^{\text{B},\text{B}2}\mathcal{B}_i$ for the individual reconstruction modes are listed in Appendix B. The values of the aforementioned quantities, as obtained from the fit, and the calculated branching fractions are summarized in Table I.

The following sources of systematic uncertainties are the most important ones and have been considered in this work: the detection-efficiency (DE) uncertainties (σ_{DE}), the statistical uncertainty on the efficiency determined from simulation (σ_{eff}), the uncertainties of the branching fractions of intermediate states ($\sigma_{\mathcal{B}_i}$), the uncertainty of the total

number of $\Upsilon(4S)$ events ($\sigma_{N_{\Upsilon(4S)}}$), the uncertainty on the fraction of charged or neutral $B\bar{B}$ events ($\sigma_{f_{+-/00}}$), the uncertainty associated with the BCS strategy (σ_{BCS}), and uncertainties related to the fit models (σ_{fit}). We explain the sources of systematics in the following and summarize their values in Table II, where the total uncertainty (σ_{total}) is calculated by adding the individual contributions in quadrature.

The DE uncertainties include contributions from tracking efficiency, PID efficiency, and the reconstruction of photon, K_S^0 , and $\bar{\Lambda}$ candidates, which are estimated using data control samples. The tracking efficiency uncertainty depends on the particle charge, momentum, and polar angle, and ranges from 0.32% to 1.22% (0.38% to 1.43%) per track in Belle (Belle II). These uncertainties are determined from control samples of $D^{*+} \rightarrow D^0(\rightarrow K_S^0\pi^+\pi^-)\pi^+$ for Belle and $\bar{B}^0 \rightarrow D^{*+}(\rightarrow D^0\pi^+)\pi^-$ and $e^+e^- \rightarrow \tau^+\tau^-$ for Belle II. The slightly larger tracking uncertainties in Belle II result from the limited statistics of the control samples. The PID efficiency uncertainties are estimated to be 1.0% (0.8%) for pions, 1.3% (1.1%) for kaons, and 2.2% (0.7%) for protons in Belle (Belle II) [30,31]. The photon reconstruction uncertainty is 2.0% (1.6%) in Belle (Belle II), as determined from control samples of radiative Bhabha (radiative muon-pair) events. The K_S^0 reconstruction uncertainty is 1.6% and 2.1% in Belle and Belle II, respectively, evaluated using control samples of $D^{*+} \rightarrow D^0(\rightarrow K_S^0\pi^0)\pi^+$ and $D^{*+} \rightarrow D^0(\rightarrow K_S^0\pi^+\pi^-)\pi^+$. The $\bar{\Lambda}$ reconstruction uncertainty is 2.2% and 1.3% in Belle and Belle II, respectively, evaluated using control samples of $\bar{\Lambda} \rightarrow \bar{p}\pi^+$ and $\bar{\Lambda}_c^- \rightarrow \bar{\Lambda}(\rightarrow \bar{p}\pi^+)\pi^-$. The individual uncertainties for the different reconstruction modes in Belle and Belle II are weighted by $N_{\Upsilon(4S)}^{\text{B},\text{B}2} \times (\epsilon_i^{\text{B},\text{B}2}\mathcal{B}_i)$ and then combined. Assuming these uncertainties are independent and adding

TABLE II. Summary of fractional systematic uncertainties (%) on the $\mathcal{B}(B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c'^-)$ and $\mathcal{B}(B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c'^0)$. The sources of all uncertainties are described in the text.

	σ_{DE}	σ_{eff}	$\sigma_{\mathcal{B}_i}$	$\sigma_{N_{\Upsilon(4S)}}$	$\sigma_{f_{+-/00}}$	σ_{BCS}	σ_{fit}	σ_{total}
B^+	2.7	1.0	3.2	1.0	2.1	2.5	4.3	7.0
B^0	2.4	1.0	3.2	1.0	1.7	3.6	4.9	7.6

them in quadrature, the DE uncertainties are evaluated to be 2.7% and 2.4% for the charged and neutral channels, respectively.

The statistical uncertainty on the simulation-based efficiency is at most 1.0%. The uncertainties on the branching fractions of $\Lambda_c^+ \rightarrow pK^-\pi^+$, $\Lambda_c^+ \rightarrow pK_S^0$, $K_S^0 \rightarrow \pi^+\pi^-$, $\bar{\Xi}_c^- \rightarrow \bar{\Xi}^+\pi^-\pi^-$, $\bar{\Xi}_c^0 \rightarrow \bar{\Xi}^+\pi^-$, $\bar{\Xi}_c^0 \rightarrow \bar{\Lambda}K^+\pi^-$, $\bar{\Xi}^+ \rightarrow \bar{\Lambda}\pi^+$, and $\bar{\Lambda} \rightarrow \bar{p}\pi^+$, and on the ratio $R = \mathcal{B}(\bar{\Xi}_c^- \rightarrow \bar{p}K^+\pi^-)/\mathcal{B}(\bar{\Xi}_c^- \rightarrow \bar{\Xi}^+\pi^-\pi^-)$, are taken from Ref. [10]. The corresponding systematic uncertainties are evaluated by varying each quantity by $\pm 1\sigma$ and taking the resulting deviation of the measured branching fraction from its nominal value as the corresponding uncertainty. For the two Λ_c^+ decay modes and the two $\bar{\Xi}_c^0$ decay modes, the correlations between the corresponding branching fractions are taken into account using the correlation coefficients provided in Ref. [10]. Correlation coefficients of 32% and 74% are used for the Λ_c^+ and $\bar{\Xi}_c^0$ decay modes, respectively. For the $\bar{\Xi}_c^-$ decay modes, the uncertainties on R and $\mathcal{B}(\bar{\Xi}_c^- \rightarrow \bar{\Xi}^+\pi^-\pi^-)$ are propagated independently. Given the large systematic uncertainties arising from the branching fractions of $\bar{\Xi}_c^-$ and $\bar{\Xi}_c^0$ decays, which may be reduced by future measurements, their contributions are quoted separately as a third uncertainty. The resulting relative systematic uncertainties due to the uncertainties on the branching fractions of $\bar{\Xi}_c^-$, $\bar{\Xi}_c^0$, and the other intermediate states are $^{+82\%}_{-33\%}$, $^{+20\%}_{-17\%}$, and 3.2%, respectively. The uncertainties of $N_{\Upsilon(4S)}^{\text{B,B2}}$ are 1.4% for Belle and 1.5% for Belle II and are combined into a total uncertainty weighted by $N_{\Upsilon(4S)}^{\text{B,B2}} \times \sum_i (\varepsilon_i^{\text{B,B2}} \mathcal{B}_i)$. The uncertainties of f_{+-} and f_{00} are 2.1% and 1.7%, respectively. To evaluate the uncertainty associated with the BCS strategy, we adopt an alternative BCS procedure in which the candidate with the minimum $|M(\gamma\bar{\Xi}_c^{\prime-0}) - m(\bar{\Xi}_c^{\prime-0})|$ is retained instead of the nominal criterion, where $m(\bar{\Xi}_c^{\prime-0})$ denotes the known mass of the $\bar{\Xi}_c^{\prime-0}$ baryon [10]. The resulting change in the measured branching fraction is taken as the corresponding systematic uncertainty. This uncertainty is estimated to be 2.5% and 3.6% for the charged and neutral channels, respectively.

The systematic uncertainties associated with the fit models arise from the signal PDF, the background PDFs, the choice of $M(\gamma\bar{\Xi}_c^{\prime-0})$ sideband regions, the modeling of the SCF components, the parameters f_{sig} and f_{scf} , and the possible peaking background from nonresonant $\Lambda_c^+\pi^\pm$ combinations in the $\Sigma_c(2455)^{++0}$ signal region. The signal PDF is replaced by a triple-Gaussian function, and the resulting systematic uncertainty is negligible. To evaluate the systematic uncertainty associated with the background PDFs, pseudoexperiments are generated using a second-order polynomial or an exponential function, with parameters taken from the corresponding alternative fits to data, while keeping the signal component unchanged. Each pseudoexperiment is fitted with the nominal background

model, and the largest difference between the mean fitted signal yield and the corresponding input yield is taken as systematic uncertainty. To assess the uncertainty associated with the choice of $M(\gamma\bar{\Xi}_c^{\prime-0})$ sideband regions, the lower and upper sideband regions are shifted simultaneously toward lower and higher mass, respectively, by 1, 2, 3, 4, and 5 MeV, with their widths kept unchanged. For the largest variation, the mass range between the lowest and highest sideband boundaries corresponds to about 15σ of the $M(\gamma\bar{\Xi}_c^{\prime-0})$ resolution. The largest deviation of the fitted signal yield from the nominal result is assigned as the systematic uncertainty. The SCF-related uncertainty is evaluated by interchanging the SCF PDFs derived from the $M(\gamma\bar{\Xi}_c^{\prime-0})$ signal and sideband regions, with the resulting change in the fitted signal yield taken as the systematic uncertainty. To estimate the systematic uncertainty associated with f_{sig} , we treat f_{sig} as a free parameter and repeat the fit. The deviation from the nominal result is taken as the systematic uncertainty. The uncertainty related to f_{scf} is covered by the variations of the sideband regions and the interchange of the SCF PDFs, which probe the uncertainty in the SCF modeling. The possible peaking background from nonresonant $\Lambda_c^+\pi^\pm$ combinations to the ΔE distribution is studied using events in the $M(\Lambda_c^+\pi^\pm)$ sideband regions, and a simultaneous fit to the ΔE distributions in the $M(\gamma\bar{\Xi}_c^{\prime-0})$ signal region, the $M(\gamma\bar{\Xi}_c^{\prime-0})$ sideband regions, and the $M(\Lambda_c^+\pi^\pm)$ sideband regions for Belle and Belle II datasets is performed. The resulting change in the fitted signal yield is smaller than 0.5% for both the charged and neutral channels and is therefore neglected. The systematic uncertainties associated with the background PDFs, the choice of $M(\gamma\bar{\Xi}_c^{\prime-0})$ sideband regions, the SCF components, and the value of f_{sig} are estimated to be 2.6% (3.1%), 2.4% (1.6%), 1.5% (2.8%), and 2.0% (2.0%) for the charged channel (neutral), respectively. All of these contributions are summed in quadrature to obtain the total systematic uncertainty related to the fit models.

In summary, we report the first observation of the decays $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime-}$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime0}$, using data samples corresponding to 771.6×10^6 and 520.6×10^6 $\Upsilon(4S)$ events collected by the Belle and Belle II experiments, respectively. This represents the first observation of B -meson decays into a pair of charmed baryon-antibaryon states that belong to the same $SU(3)$ flavor sextet. The branching fractions are measured to be $\mathcal{B}(B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime-}) = (1.69 \pm 0.31 \pm 0.12_{-0.56}^{+1.39}) \times 10^{-3}$ and $\mathcal{B}(B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime0}) = (1.28 \pm 0.32 \pm 0.10_{-0.22}^{+0.26}) \times 10^{-3}$, where the uncertainties are statistical, systematic, and from the absolute branching fractions of $\bar{\Xi}_c^-$ and $\bar{\Xi}_c^0$ decays, respectively. The measured branching fractions are 2–3 times larger than those of $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^-$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^0$ [11], despite the smaller available phase space for these decays. This observation may

indicate nontrivial dynamical effects in baryonic B decays that enhance or suppress specific decay modes.

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DATA AVAILABILITY

There are no publicly available research data or software supporting this manuscript. Requests for further information or data should be sent to the authors.

APPENDIX A: $M(\Lambda_c^+\pi^\pm)$ VERSUS $M(\gamma\bar{\Xi}_c^{\prime -0})$ DISTRIBUTIONS

Figure 3 shows the scatter plots of $M(\Lambda_c^+\pi^\pm)$ versus $M(\gamma\bar{\Xi}_c^{\prime -0})$ for reconstructed $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime -}$ and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime 0}$ candidates in the signal and sideband regions of ΔE from the combined Belle and Belle II data samples, before application of the BCS. The signal region of ΔE is defined as $|\Delta E| < 0.012$ GeV (approximately 2.5σ), while the sideband regions are defined as $0.030 < |\Delta E| < 0.054$ GeV, with a total width twice that of the signal region. The corresponding one-dimensional projections onto $M(\Lambda_c^+\pi^\pm)$ and $M(\gamma\bar{\Xi}_c^{\prime -0})$ for candidates in the signal region of ΔE are shown in Fig. 4. Clear signals corresponding to $\Sigma_c(2455)^{+,0}$ and $\bar{\Xi}_c^{\prime -,0}$ are observed.

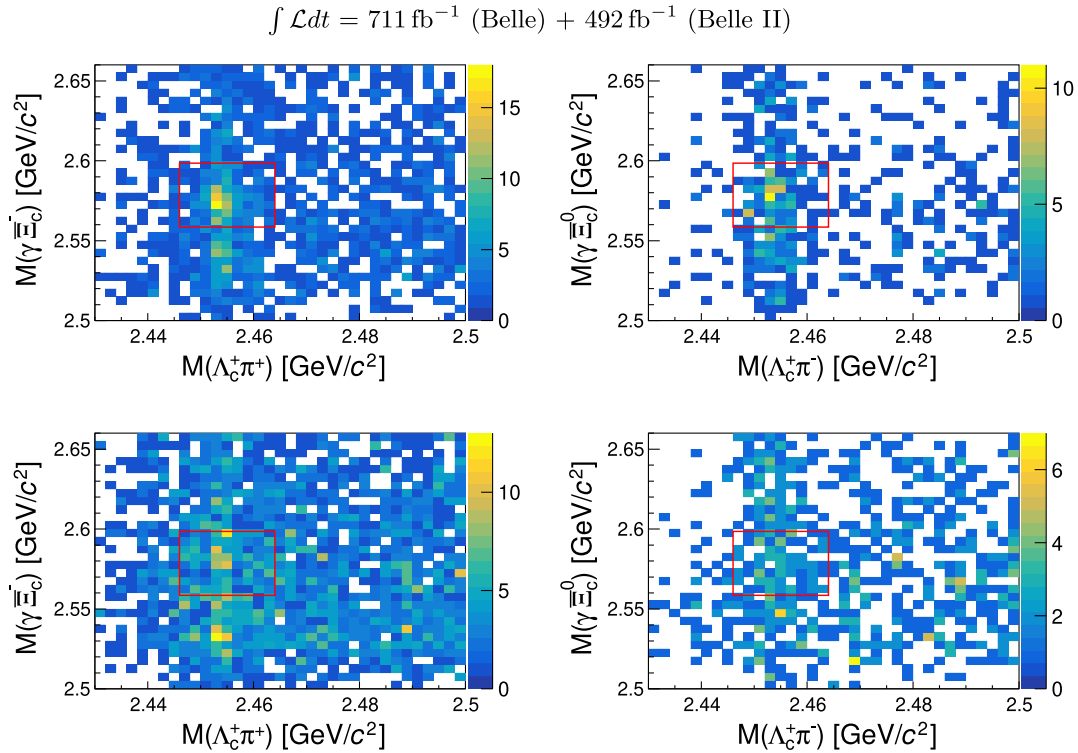


FIG. 3. Scatter plots of $M(\Lambda_c^+\pi^\pm)$ versus $M(\gamma\bar{\Xi}_c^{\prime -0})$ for reconstructed $B^+ \rightarrow \Sigma_c(2455)^{++}\bar{\Xi}_c^{\prime -}$ (left) and $B^0 \rightarrow \Sigma_c(2455)^0\bar{\Xi}_c^{\prime 0}$ (right) candidates in the signal (top) and sideband (bottom) regions of ΔE from the combined Belle and Belle II data samples, before application of the BCS. The red boxes indicate the signal regions for the $\Sigma_c(2455)^{+,0}$ and $\bar{\Xi}_c^{\prime -,0}$ candidates.

$$\int \mathcal{L} dt = 711 \text{ fb}^{-1} \text{ (Belle)} + 492 \text{ fb}^{-1} \text{ (Belle II)}$$

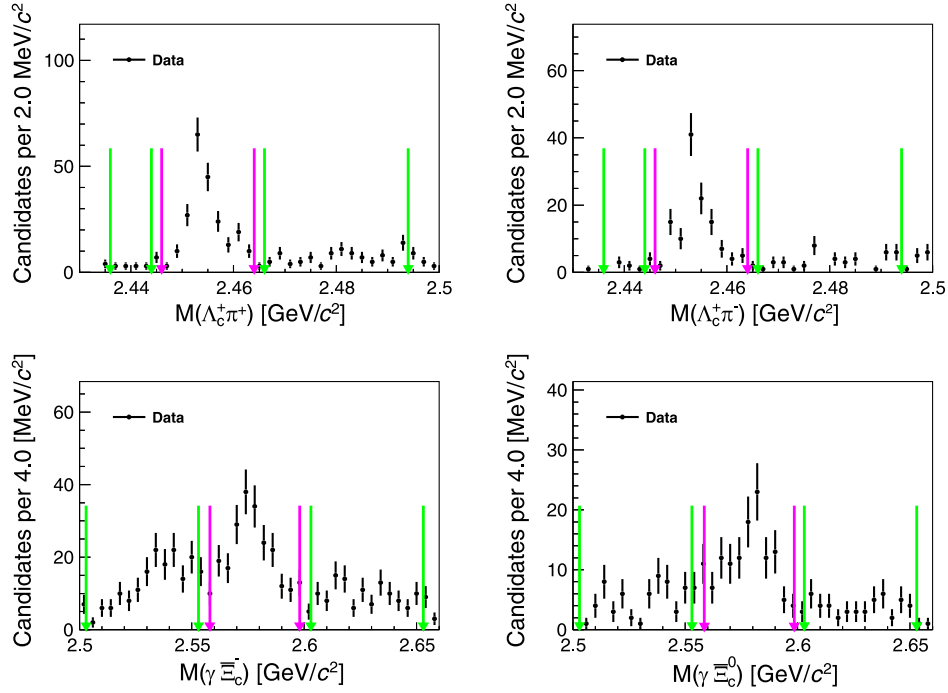


FIG. 4. One-dimensional projections onto (top) $M(\Lambda_c^+ \pi^+)$ and (bottom) $M(\gamma \bar{\Xi}_c'^0)$ for reconstructed $B^+ \rightarrow \Sigma_c(2455)^{++} \bar{\Xi}_c'^{-}$ (left) and $B^0 \rightarrow \Sigma_c(2455)^0 \bar{\Xi}_c'^0$ (right) candidates in the signal region of ΔE from combined Belle and Belle II data samples, before application of the BCS. Points with error bars represent the data. The magenta arrows indicate the signal regions defined for the $\Sigma_c(2455)^{++0}$ and $\bar{\Xi}_c'^{0}$ candidates, while the green arrows indicate the corresponding sideband regions.

APPENDIX B: RECONSTRUCTION EFFICIENCIES FOR THE INDIVIDUAL RECONSTRUCTION MODES

Tables III and IV list the reconstruction efficiencies and the products of the reconstruction efficiencies and the corresponding branching fractions of all intermediate states for the individual reconstruction modes of the $B^+ \rightarrow \Sigma_c(2455)^{++} \bar{\Xi}_c'^{-}$ and $B^0 \rightarrow \Sigma_c(2455)^0 \bar{\Xi}_c'^0$ decays, respectively.

TABLE III. Reconstruction efficiencies (ϵ_i) and the products of the reconstruction efficiencies and the branching fractions of all intermediate states ($\epsilon_i \mathcal{B}_i$) for the individual reconstruction modes of the $B^+ \rightarrow \Sigma_c(2455)^{++} (\rightarrow \Lambda_c^+ \pi^+) \bar{\Xi}_c'^{-} (\rightarrow \gamma \bar{\Xi}_c^0)$ decay. The branching fractions of all intermediate states are taken from Ref. [10]. The branching fraction $\mathcal{B}(\bar{\Xi}_c^0 \rightarrow \bar{p} K^+ \pi^-)$ is inferred from the ratio $\mathcal{B}(\bar{\Xi}_c^0 \rightarrow \bar{p} K^+ \pi^-) / \mathcal{B}(\bar{\Xi}_c^0 \rightarrow \bar{\Xi}^+ \pi^- \pi^-)$ and the absolute branching fraction $\mathcal{B}(\bar{\Xi}_c^0 \rightarrow \bar{\Xi}^+ \pi^- \pi^-)$.

Reconstruction mode	Belle		Belle II	
	ϵ_i (%)	$\epsilon_i \mathcal{B}_i$ (10^{-5})	ϵ_i (%)	$\epsilon_i \mathcal{B}_i$ (10^{-5})
$\Lambda_c^+ \rightarrow p K^- \pi^+, \bar{\Xi}_c'^- \rightarrow \bar{\Xi}^+ \pi^- \pi^-$	0.70	0.83	1.04	1.23
$\Lambda_c^+ \rightarrow p K^- \pi^+, \bar{\Xi}_c'^- \rightarrow \bar{p} K^+ \pi^-$	3.06	1.19	3.98	1.54
$\Lambda_c^+ \rightarrow p K_S^0, \bar{\Xi}_c'^- \rightarrow \bar{\Xi}^+ \pi^- \pi^-$	0.82	0.17	1.17	0.25
$\Lambda_c^+ \rightarrow p K_S^0, \bar{\Xi}_c'^- \rightarrow \bar{p} K^+ \pi^-$	3.43	0.24	4.40	0.30
$\sum_i \epsilon_i \mathcal{B}_i$	2.43×10^{-5}		3.32×10^{-5}	

TABLE IV. Reconstruction efficiencies (ϵ_i) and the products of the reconstruction efficiencies and the branching fractions of all intermediate states ($\epsilon_i \mathcal{B}_i$) for the individual reconstruction modes of the $B^0 \rightarrow \Sigma_c(2455)^0 (\rightarrow \Lambda_c^+ \pi^-) \bar{\Xi}_c^{\prime 0} (\rightarrow \gamma \bar{\Xi}_c^0)$ decay. The branching fractions of all intermediate states are taken from Ref. [10].

Reconstruction mode	Belle		Belle II	
	ϵ_i (%)	$\epsilon_i \mathcal{B}_i$ (10^{-5})	ϵ_i (%)	$\epsilon_i \mathcal{B}_i$ (10^{-5})
$\Lambda_c^+ \rightarrow p K^- \pi^+, \bar{\Xi}_c^0 \rightarrow \bar{\Xi}^+ \pi^-$	1.04	0.61	1.46	0.85
$\Lambda_c^+ \rightarrow p K^- \pi^+, \bar{\Xi}_c^0 \rightarrow \bar{\Lambda} K^+ \pi^-$	1.13	0.67	2.05	1.22
$\Lambda_c^+ \rightarrow p K_S^0, \bar{\Xi}_c^0 \rightarrow \bar{\Xi}^+ \pi^-$	1.18	0.12	1.67	0.17
$\Lambda_c^+ \rightarrow p K_S^0, \bar{\Xi}_c^0 \rightarrow \bar{\Lambda} K^+ \pi^-$	1.28	0.13	2.33	0.25
$\sum_i \epsilon_i \mathcal{B}_i$	1.53×10^{-5}		2.49×10^{-5}	

- [1] H. Y. Cheng, *Nucl. Phys. B, Proc. Suppl.* **163**, 68 (2007).
- [2] R. Chistov, *J. Phys. Conf. Ser.* **675**, 022007 (2016).
- [3] Ed. A. J. Bevan, B. Golob, Th. Mannel, S. Prell, and B. D. Yabsley, *Eur. Phys. J. C* **74**, 3026 (2014).
- [4] Y. K. Hsiao, S. Y. Tsai, C. C. Lih, and E. Rodrigues, *J. High Energy Phys.* **04** (2020) 035.
- [5] Y. K. Hsiao, *J. High Energy Phys.* **11** (2023) 117.
- [6] R. Chistov *et al.* (Belle Collaboration), *Phys. Rev. D* **74**, 111105 (2006).
- [7] B. Aubert *et al.* (BABAR Collaboration), *Phys. Rev. D* **77**, 031101 (2008).
- [8] Y. Li *et al.* (Belle Collaboration), *Phys. Rev. D* **100**, 112010 (2019).
- [9] R. Aaij *et al.* (LHCb Collaboration), *Phys. Rev. Lett.* **136**, 061802 (2026).
- [10] F. Takahashi *et al.* (Particle Data Group), *Int. J. Mod. Phys. A* **41**, 2630011 (2026).
- [11] M. Abumusabh *et al.* (Belle Collaboration and Belle II Collaboration), *Phys. Rev. D* **112**, L051101 (2025).
- [12] A. Abashian *et al.* (Belle Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **479**, 117 (2002).
- [13] T. Abe *et al.* (Belle II Collaboration), *arXiv:1011.0352*.
- [14] J. Brodzicka *et al.* (Belle Collaboration), *Prog. Theor. Exp. Phys.* **2012**, 04D001 (2012).
- [15] I. Adachi *et al.* (Belle Collaboration), *Chin. Phys. C* **49**, 013001 (2025).
- [16] S. Kurokawa and E. Kikutani, *Nucl. Instrum. Methods Phys. Res., Sect. A* **499**, 1 (2003); T. Abe *et al.*, *Prog. Theor. Exp. Phys.* **2013**, 03A001 (2013).
- [17] K. Akai *et al.* (SuperKEKB Accelerator Team), *Nucl. Instrum. Methods Phys. Res., Sect. A* **907**, 188 (2018).
- [18] D. J. Lange, *Nucl. Instrum. Methods Phys. Res., Sect. A* **462**, 152 (2001).
- [19] T. Sjöstrand, S. Mrenna, and P. Z. Skands, *J. High Energy Phys.* **05** (2006) 026.
- [20] T. Sjöstrand *et al.*, *Comput. Phys. Commun.* **191**, 159 (2015).
- [21] E. Barberio, B. van Eijk, and Z. Was, *Comput. Phys. Commun.* **66**, 115 (1991).
- [22] Z. J. Liptak *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. A* **1040**, 167168 (2022).
- [23] X. Y. Zhou, S. X. Du, G. Li, and C. P. Shen, *Comput. Phys. Commun.* **258**, 107540 (2021).
- [24] S. Jadach, B. F. L. Ward, and Z. Was, *Comput. Phys. Commun.* **130**, 260 (2000).
- [25] R. Brun *et al.*, GEANT3: user's guide GEANT3.10, GEANT3.11, CERN Report No. DD/EE/84-1, 1984.
- [26] S. Agostinelli *et al.* (GEANT4 Collaboration), *Nucl. Instrum. Methods Phys. Res., Sect. A* **506**, 250 (2003).
- [27] T. Kuhr, C. Pulvermacher, M. Ritter, T. Hauth, and N. Braun (Belle II Framework Software Group), *Comput. Software Big Sci.* **3**, 1 (2019).
- [28] M. Gelb *et al.*, *Comput. Software Big Sci.* **2**, 9 (2018).
- [29] G. Punzi, *eConf C030908*, MODT002 (2003).
- [30] E. Nakano, *Nucl. Instrum. Methods Phys. Res., Sect. A* **494**, 402 (2002).
- [31] I. Adachi *et al.* (Belle II Collaboration), *Eur. Phys. J. C* **85**, 1237 (2025).
- [32] H. Nakano *et al.* (Belle Collaboration), *Phys. Rev. D* **97**, 092003 (2018).
- [33] F. Abudinén *et al.* (Belle Collaboration and Belle II Collaboration), *J. High Energy Phys.* **02** (2022) 063.
- [34] J. Y. Lee *et al.* (Belle Collaboration), *Phys. Rev. D* **103**, 052005 (2021).
- [35] P. Cheema, *EPJ Web Conf.* **295**, 09035 (2024).
- [36] K. S. Cranmer, *Comput. Phys. Commun.* **136**, 198 (2001).
- [37] S. Banerjee *et al.* (Heavy Flavor Averaging Group), *Phys. Rev. D* **113**, 012008 (2026), with specific result from.