

Long-lived axionlike particles from tau decays

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(Received 12 August 2025; accepted 17 November 2025; published 15 December 2025)

Axionlike particles (ALPs) are well-motivated examples of light, weakly coupled particles in theories beyond the Standard Model. In this work, we study long-lived ALPs coupled exclusively to leptons in the mass range between $2m_e$ and $m_\tau - m_e$. For anarchic flavor structure the leptophilic ALP production in tau decays or from ALP-tau bremsstrahlung is enhanced thanks to derivative couplings of the ALP and can surpass production from electron and muon channels, especially for ALPs heavier than m_μ . Using past data from high-energy fixed-target experiments such as CHARM and BEBC we place new constraints on the ALP decay constant f_a , reaching scales as high as $\mathcal{O}(10^8)$ GeV in lepton-flavor-violating channels and $f_a \sim \mathcal{O}(10^2)$ GeV in lepton-flavor-conserving ones. We also present projections for the event-rate sensitivity of current and future detectors to ALPs produced at the Fermilab Main Injector, the CERN SPS, and in the forward direction of the LHC. We show that SHiP will be sensitive to f_a values that are over an order of magnitude above the existing constraints.

DOI: [10.1103/51gx-m32t](https://doi.org/10.1103/51gx-m32t)

I. INTRODUCTION

The search for light, very weakly coupled particles has emerged as a crucial avenue for exploring physics beyond the Standard Model (SM). This approach is complementary to the more traditional strategy of searching for new heavy particles at increasingly higher energies. Over the last two decades, significant efforts have been devoted to defining a “space of possibilities” for feebly interacting particles, commonly referred to as the “dark sector physics” (see, e.g., Refs. [1–3] and references therein).

One of the better motivated examples of such physics is a light scalar, a , derivatively coupled to the SM fermions $\psi_{L,R}^i$,

$$\mathcal{L}_{\text{int}} \propto (\bar{\psi}_{L,R}^i \gamma^\mu \psi_{L,R}^j) \times \partial_\mu a / f_a. \quad (1.1)$$

These interactions are of canonical dimension 5, and are thus suppressed by an energy scale, conventionally denoted as f_a . Similar types of derivatively coupled operators can be built from the SM gauge fields,

$$(\epsilon^{\mu\nu\alpha\beta} \text{Tr}[A_\nu \partial_\alpha A_\beta]) \times \partial_\mu a / f_a, \quad (1.2)$$

with $A_\mu = A_\mu^a T^a$ denoting any of the $SU(3)_c \times SU(2)_L \times U(1)_Y$ gauge fields. A notable feature of these interactions is the absence of large perturbative corrections to the a mass, m_a , from loops involving the SM fermions $\psi_{L,R}^i$. At a nonperturbative level, however, the interactions with gluons do generate a mass term for a , which plays a crucial role in the QCD axion mechanism—allowing the originally massless axion field to dynamically restore CP conservation in QCD [4–6].

In general, a light scalar a need not acquire its mass solely through the QCD anomaly. This a more general

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possibility, where m_a is treated as a free parameter, it will be referred to as an axionlike particle (ALP). It is also worth noting that ALP interactions with the SM fermions constitute only a subset of a broader class of operators, which include also nonderivative and/or scalar couplings, as recently discussed in [7,8].

In the general interaction of Eq. (1.1), the indices i, j label different types of SM fermion, such as quarks and leptons and include different generations (respecting, of course, an overall gauge invariance of the current). Clearly, a very large variety of such couplings is possible, and without the knowledge of ultraviolet completion of these interactions no definitive predictions regarding their relative size can be made. In this paper, we will stay at the level of the effective field theories (EFTs), and will explore couplings of a to leptonic currents. Our primary emphasis will be on the relatively unexplored class of interactions involving tau leptons. The lowest dimension EFT Lagrangian is thus

$$\mathcal{L}_a = \mathcal{L}_{\text{kin}} + \mathcal{L}_{\text{int}}, \quad (1.3)$$

where the canonical kinetic term for a ,

$$\mathcal{L}_{\text{kin}} = \frac{1}{2}(\partial_\mu a)^2 - \frac{1}{2}m_a^2 a^2, \quad (1.4)$$

is supplemented by the interaction terms involving just the lepton currents,

$$\mathcal{L}_{\text{int}} = \frac{\partial_\mu a}{2f_a} \sum_{i,j} [g_L^{ij} \bar{L}_i \gamma^\mu P_L L_j + g_R^{ij} \bar{\ell}_i \gamma^\mu P_R \ell_j], \quad (1.5)$$

where we made the chirality explicit and write the interaction in the $\text{SU}(2)_L$ -symmetric phase with $L_i = (\nu_i)$ the lepton doublet. In this expression, f_a is the overall suppression scale, while the indices on the dimensionless real coefficients g_L^{ij} and g_R^{ij} run over the three generations of SM leptons, $i, j = 1, 2, 3$.

In the analysis and numerical calculation below we limit ourselves to the situation where g_L^{ij} and g_R^{ij} have the same magnitude but opposite sign, i.e., $g_{\ell_1 \ell_2} = g_L^{\ell_1 \ell_2} = -g_R^{\ell_1 \ell_2}$, and use the following short-hand notation,

$$\frac{1}{f_{\ell_1 \ell_2}} = \frac{g_{\ell_1 \ell_2}}{f_a}. \quad (1.6)$$

The Lagrangian in Eq. (1.5) has two phenomenologically important features. The off-diagonal couplings $f_{e\tau}^{-1}$ and $f_{\mu\tau}^{-1}$ induce flavor-changing neutral current (FCNC) decays $\tau \rightarrow e a$ and $\tau \rightarrow \mu a$. Since these are induced by dimension-5 operators they are parametrically enhanced by $(m_W^2/m_\tau f_{\ell\tau})^2$ compared to the SM decay rates, which are suppressed by Fermi coupling constant $G_F \propto 1/m_W^2$. Therefore, nonzero

$f_{\mu\tau}^{-1}$ and $f_{e\tau}^{-1}$ can efficiently produce a in any process that results in τ leptons (see, for example, Ref. [9]).

The second notable feature of the Lagrangian in Eq. (1.5) is the additional suppression of the partial decay widths $a \rightarrow \mu^+ \mu^-$ and, in particular, $a \rightarrow e^+ e^-$. The currents that the axion couples to are conserved in the limit of massless leptons, so the two decays widths are proportional to lepton masses. In particular, $\Gamma(a \rightarrow \mu^+ \mu^-) \propto m_\mu (m_\mu/f_{\mu\mu})^2$ and $\Gamma(a \rightarrow e^+ e^-) \propto m_e (m_e/f_{ee})^2$, and are thus suppressed by the relative smallness of the leptonic masses m_μ and m_e relative to the fiducial mass scale set by m_a , which can be as large as ~ 1 GeV. As a result, a is expected to be relatively long-lived, potentially giving rise to displaced decay signatures—a possibility that we explore in this work.

Even though we restricted the parameter space of possible ALP couplings to just the ones coupling ALPs to leptons, there is still considerable freedom in the choices for these couplings. At the qualitative level, though, the bulk of the phenomenology can be captured by considering the following lepton-flavor-violating (LFV) and lepton-flavor-conserving (LFC) benchmark scenarios:

- (1) Anarchical LFV. This scenario assumes a “democratic” pattern of couplings, without any particular suppression of the off-diagonal couplings,

$$g_{\tau\tau} \sim g_{\mu\tau} \sim g_{\mu\mu} \sim g_{ee}, \text{ etc.} \quad (1.7)$$

The main source of a will be the FCNC production from $\tau \rightarrow \ell a$, while the dominant a decay channel will be into the heaviest kinematically available lepton pair.

- (2) Hierarchical LFV. In this scenario the flavor-changing couplings are assumed to be smaller,

$$g_{\tau\tau} \sim g_{\mu\mu} \sim g_{ee} > g_{\mu\tau}, g_{e\tau}, g_{e\mu}, \quad (1.8)$$

but not negligible, so that the FCNC decays $\tau \rightarrow \ell a$ still dominate ALP production. Within this scenario, we will be able to explore how the size of the hierarchy between LFV and LFC couplings impacts the experimental reach.

- (3) Universal LFC. In this benchmark, the FCNC couplings are set to zero, while the diagonal couplings have no particular pattern

$$g_{\tau\tau} \sim g_{\mu\mu} \sim g_{ee}; \quad g_{\mu\tau} = g_{e\tau} = g_{e\mu} = 0. \quad (1.9)$$

In this scenario, the “bremsstrahlung” production of a in tau decays (e.g., $\tau \rightarrow \pi \nu a; \pi \pi \nu a$ etc) will be important, and the longevity of a will be achieved only for $m_a < 2m_\mu$.

- (4) Nonuniversal LFC. In this benchmark, ALPs only have flavor diagonal couplings to leptons, and these exhibit a hierarchical pattern. We are particularly interested in the case, where ALPs produced in tau

decays are always long-lived, which occurs for a flavor nonuniversal coupling choice of

$$g_{\mu\mu} \ll g_{\tau\tau}; \quad g_{\mu\tau} = g_{e\tau} = g_{e\mu} = 0. \quad (1.10)$$

Depending on whether g_{ee} is sizeable or not, the ALP will decay mostly to e^+e^- or to $\gamma\gamma$, but it will be typically long-lived in the mass range of interest. We consider both cases: a μ -*phobic* ALP, with

$$g_{\mu\mu} \ll g_{ee} \sim g_{\tau\tau}, \quad (1.11)$$

and a τ -*philic* scenario with

$$g_{\mu\mu}, g_{ee} \ll g_{\tau\tau}. \quad (1.12)$$

Our goal is to obtain existing constraints on the above scenarios and provide projections both for the upgrades of the existing experiments, as well as for future projects such as SHiP, focusing on processes involving tau leptons. The main finding is that beam dump and collider searches for displaced ALPs from tau decays provide better sensitivity to such ALPs than the precision studies of tau decays. The displaced ALP decays to $\mu^+\mu^-$ and muon-electron pairs are a currently underexplored signature of LFV and LFC models of ALPs coupled exclusively to lepton currents, even though both channels are generically available and even dominant. We make projections for these types of searches for SHiP and the next generation of the FASER experiment. In the LFC case, we also derive novel constraints on light ALPs with displaced decays to e^+e^- pairs.

The paper is organized as follows. In Sec. II we derive the branching ratios for production of ALPs in τ decays, as well as ALP decay rates. In Sec. III we discuss the resulting ALP production rates at past, existing, and future experiments. Our main results, the current and projected limits on the ALP parameter space, are given in Sec. IV. Section V contains our conclusions, while further details on the calculation of ALP production and decay rates are relegated to Appendix C.

II. ALP PRODUCTION AND DECAY

In this section, we summarize the ALP production and decay rates. We focus on the production of ALPs in the decays of tau leptons, see Fig. 1. The corresponding differential decay rate distributions, in τ rest frame, are shown in Fig. 2. Additional details are provided in Appendix C. The main production mechanism of τ leptons in beam dump experiments are the decays of D_s meson, since the branching ratio $\text{Br}(D_s^\pm \rightarrow \tau^\pm \nu_\tau^{(-)}) \simeq 5.36\%$ is quite large. Other subleading production channels include $\text{Br}(D^\pm \rightarrow \tau^\pm \nu_\tau^{(-)}) = 1.20 \times 10^{-3}$ and $\text{Br}(\psi(2S) \rightarrow \tau^+ \tau^-) = 3.10 \times 10^{-3}$ decays [10]. The dominant production of ALPs is then either from

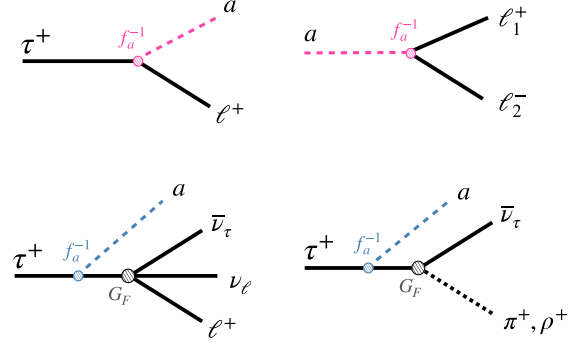


FIG. 1. Tau decays involving LFV (top left) and LFC (bottom row) ALPs, where, in addition, the ALP can also attach to $\bar{\nu}_\tau$ (not shown). The ALP decay is shown in the top right.

two-body decays of tau leptons (for the LFV scenarios) or from multibody decays of taus (for the LFC scenarios).

A. ALP production: LFV scenarios

If the LFV couplings of the ALP are appreciable, the ALP is dominantly produced through the flavor-changing decay of tau leptons, $\tau \rightarrow \mu a, e a$. This is the case for both the anarchical and hierarchical LFV scenarios, Eqs. (1.7) and (1.8), respectively. The decay rate is given by

$$\Gamma(\tau \rightarrow \ell a) = \frac{|g_{\ell\tau}|^2 m_\tau^3}{64\pi f_a^2} \left(1 + \frac{m_\ell}{m_\tau}\right)^2 \left[\left(1 - \frac{m_\ell}{m_\tau}\right)^2 - \frac{m_a^2}{m_\tau^2} \right] \times v(m_\tau, m_\ell, m_a), \quad (2.1)$$

where $\ell = \mu, e$, and v is the separation velocity,

$$v(M, m_1, m_2) = \sqrt{1 - \frac{2(m_1^2 + m_2^2)}{M^2} + \frac{(m_1^2 - m_2^2)^2}{M^4}}. \quad (2.2)$$

For instance, for $m_a = 0.1$ GeV, we find

$$\text{Br}(\tau \rightarrow \ell a) = 1.2 \times 10^{-4} |g_{\ell\tau}|^2 \left(\frac{10^7 \text{ GeV}}{f_a} \right)^2, \quad (2.3)$$

with $\ell = e, \mu$. Note that the branching ratio for ALP production is enhanced by the small SM tau decay width, i.e., it is enhanced by G_F^{-2} , where G_F is the Fermi constant. Consequently, $\text{Br}(\tau \rightarrow \ell a)$ is sizable even for relatively large values of f_a , as long as $g_{\ell\tau}$ is not too small. A large value of f_a , on the other hand, implies a long-lived ALP.

Here we note that, if $m_a < m_\mu - m_e$, the $\mu \rightarrow ea$ decays are kinematically allowed. In this ALP mass range the searches for $\mu \rightarrow ea$ decays place stringent constraints on ALP LFV scenarios as long as $g_{e\mu}$ coupling is comparable to the $g_{\ell\tau}$ couplings [11–14].

B. ALP production: LFC scenarios

If flavor-changing couplings $g_{\ell\tau}$ are vanishingly small, i.e., for the LFC scenarios in Eqs. (1.9) and (1.10), tau leptons can still produce ALPs. The ALP production then proceeds via ALP bremsstrahlung modifying any of the SM τ decay channels, see the bottom row in Fig. 1. For simplicity, we focus on the decay modes $\tau^- \rightarrow \pi^- \nu_\tau$, $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$, and $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau$, which account for $\sim 70\%$ of the SM τ decay width. Including more decay modes would result in a marginally tighter bound on (better sensitivity to) the ALP couplings.

To calculate the ALP production via bremsstrahlung in $\tau \rightarrow \pi \nu_\tau a$, we use the following effective Lagrangian for the SM rate

$$\mathcal{L}_{\text{eff}} = -\frac{G_F}{\sqrt{2}} V_{ud}^* f_\pi (\partial_\mu \pi^-) \bar{\tau} \gamma^\mu (1 - \gamma_5) \nu_\tau + (\text{H.c.}), \quad (2.4)$$

where $V_{ud} \simeq 0.97$ is the Cabibbo–Kobayashi–Maskawa matrix element and $f_\pi = 130$ MeV the pion decay constant. The differential $\tau \rightarrow \pi \nu_\tau a$ decay rate is given by¹

$$\begin{aligned} \frac{d\Gamma(\tau \rightarrow \pi \nu_\tau a)}{dE_a} &= \frac{|g_{\tau\tau}|^2 G_F^2 |V_{ud}|^2 f_\pi^2 m_\tau^2 (\bar{m}_\tau^2 - m_\pi^2)^2}{256\pi^3 f_a^2 \bar{m}_\tau^2} \\ &\times \left[1 - \frac{m_a^2 (m_\tau^2 + \bar{m}_\tau^2)}{(m_\tau^2 - \bar{m}_\tau^2)^2} \right] v(m_\tau, \bar{m}_\tau, m_a), \end{aligned} \quad (2.5)$$

where E_a is the ALP energy in the τ -rest frame. Because of ALP bremsstrahlung from the τ leg, the intermediate τ line is off-shell, with the off-shell invariant mass squared $\bar{m}_\tau^2$ given by

$$\bar{m}_\tau^2 = m_\tau^2 + m_a^2 - 2m_\tau E_a. \quad (2.6)$$

Integrating the differential rate in Eq. (2.5) over $E_a \in [m_a, (m_\tau^2 + m_a^2 - m_\pi^2)/2m_\tau]$, gives $\Gamma(\tau \rightarrow \pi \nu_\tau a)$. For example, for $m_a = 0.1$ GeV, we find

$$\text{Br}(\tau \rightarrow \pi \nu_\tau a) = 1.5 \times 10^{-12} |g_{\tau\tau}|^2 \left(\frac{10^4 \text{ GeV}}{f_a} \right)^2. \quad (2.7)$$

Compared to the $\tau \rightarrow \ell a$ decay rate in Eq. (2.3) this branching ratio is significantly more suppressed, if all couplings of ALP to leptons are $\mathcal{O}(1)$. The reason is that the $\tau \rightarrow \pi \nu_\tau a$ decay rate suffers double suppression, the factor of G_F^2 due to SM weak interaction, and $(g_{\tau\tau}/f_a)^2$ for the emission of the ALP. For LFC scenarios the phenomenologically relevant values of f_a will be relatively low, $f_a \lesssim 10^4$ GeV. This is true for the other tau decay channels

¹Note that since ALP couplings in Eq. (1.5) are $\text{SU}(2)_L$ -invariant, there are no so-called “weak-violating” ALP type contributions to these decays [15,16].

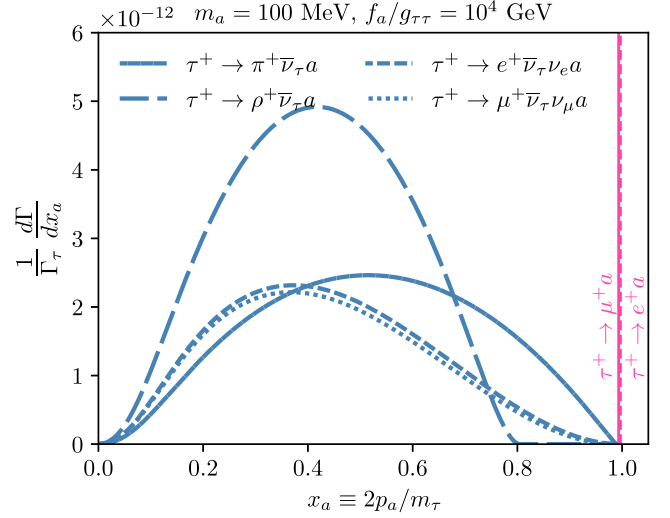


FIG. 2. Differential branching ratios for various tau decays involving ALPs, as indicated, for a sample ALP mass $m_a = 100$ MeV, and ALP decay constant $f_a = 10^4$ GeV, under the τ -philic scenario hypothesis of Eq. (1.12). Here, $x_a = 2p_a/m_\tau$ is the fractional momentum carried by the ALP in the tau rest frame (p_a is the magnitude of the ALP momentum). We also indicate with vertical purple lines the value of x_a corresponding to the two monoenergetic two-body decay modes, $\tau \rightarrow \ell a$.

involving ALPs, which we discuss below. For such low values of f_a the ALP decays are displaced only for $m_a < 2m_\mu$ or $|g_{\mu\mu}| \ll 1$, so that the dominant ALP decay is $a \rightarrow e^+ e^-$, whose rate is suppressed by m_e^2 in the assumed LFC flavor patterns.

The decay mode $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$ is dominated by the intermediate ρ meson decaying via its main decay mode, $\rho^- \rightarrow \pi^- \pi^0$. Ignoring the ρ width effects we can therefore treat this ALP production channel as though it is due to a quasi-three-body decay $\tau^- \rightarrow \rho^- \nu_\tau a$. The effective Lagrangian relevant for the ρ meson production can be written as [17]

$$\mathcal{L}_{\text{eff}} = -\frac{G_F}{\sqrt{2}} V_{ud}^* m_\rho f_\rho \rho_\mu \bar{\tau} \gamma^\mu (1 - \gamma_5) \nu_\tau + (\text{H.c.}), \quad (2.8)$$

where $m_\rho \simeq 775$ MeV is the ρ meson mass and $f_\rho = 216$ MeV the ρ decay constant. The differential decay rate is given by

$$\begin{aligned} \frac{d\Gamma(\tau \rightarrow \rho \nu_\tau a)}{dE_a} &= \frac{|g_{\tau\tau}|^2 G_F^2 |V_{ud}|^2 f_\rho^2 m_\tau^2 (\bar{m}_\tau^2 + 2m_\rho^2)}{256\pi^3 f_a^2} \\ &\times \left[1 - \frac{m_\rho^2}{\bar{m}_\tau^2} \right]^2 \left[1 - \frac{m_a^2 (m_\tau^2 + \bar{m}_\tau^2)}{(m_\tau^2 - \bar{m}_\tau^2)^2} \right] \\ &\times v(m_\tau, \bar{m}_\tau, m_a). \end{aligned} \quad (2.9)$$

Integrating the above differential rate from $E_a = m_a$ to $E_a = (m_\tau^2 + m_a^2 - m_\rho^2)/2m_\tau$ gives the $\Gamma(\tau \rightarrow \rho \nu_\tau a)$ decay rate. For example, for $m_a = 0.1$ GeV, we find

$$\text{Br}(\tau \rightarrow \rho \nu_\tau a) = 2.3 \times 10^{-12} |g_{\tau\tau}|^2 \left(\frac{10^4 \text{ GeV}}{f_a} \right)^2, \quad (2.10)$$

and it is thus numerically comparable to $\text{Br}(\tau \rightarrow \pi \nu_\tau a)$, Eq. (2.7), as long as ALP is not too heavy.

Finally, the leptonic decay mode $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau$ is governed by the four-Fermi interaction

$$\mathcal{L} = -\frac{G_F}{\sqrt{2}} [\bar{\nu}_\tau \gamma_\alpha (1 - \gamma_5) \tau] [\bar{\ell} \gamma^\alpha (1 - \gamma_5) \nu_\ell], \quad (2.11)$$

where $\ell = e, \mu$. The $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau a$ channel is a four-body decay. We perform the required phase space integrals analytically in Appendix C, leading to the following expression for the double differential decay rate

$$\begin{aligned} \frac{d\Gamma(\tau \rightarrow \ell \bar{\nu}_\ell \nu_\tau a)}{dE_a d\bar{m}_{\nu\nu}^2} &= \frac{G_F^2 |g_{\tau\tau}|^2 m_\tau^2 \bar{m}_\tau^2}{1536 \pi^5 f_a^2} I(\bar{m}_\tau, \bar{m}_{\nu\nu}) \\ &\times \left[1 - \frac{m_a^2 (m_\tau^2 + \bar{m}_\tau^2)}{(m_\tau^2 - \bar{m}_\tau^2)^2} \right] v(m_\tau, \bar{m}_\tau, m_a), \end{aligned} \quad (2.12)$$

where $E_a \in [m_a, (m_\tau^2 + m_a^2 - m_\ell^2)/2m_\tau]$ is the ALP energy, $\bar{m}_{\nu\nu}^2 \in [0, (\bar{m}_\tau - m_l)^2]$ is the invariant mass squared of the $\bar{\nu}_\ell \nu_\tau$ pair, and

$$\begin{aligned} I(\bar{m}_\tau, \bar{m}_{\nu\nu}) &= \left[1 + \frac{\bar{m}_{\nu\nu}^2 - 2m_l^2}{\bar{m}_\tau^2} + \frac{m_l^4 + m_l^2 \bar{m}_{\nu\nu}^2 - 2\bar{m}_\tau^4}{\bar{m}_\tau^4} \right] \\ &\times v(\bar{m}_\tau, \bar{m}_{\nu\nu}, m_\ell). \end{aligned} \quad (2.13)$$

In Eq. (2.12) we only included ALP emissions from τ and ν_τ lines since the ones from ℓ and $\bar{\nu}_\ell$ are suppressed by m_ℓ/m_τ at the level of the amplitude, and are thus never important for any of the LFC flavor structure we consider.

Integrating the differential rate in Eq. (2.12) over E_a and $\bar{m}_{\nu\nu}^2$ gives the $\Gamma(\tau \rightarrow \ell \bar{\nu}_\ell \nu_\tau a)$ decay rate. For instance, for $m_a = 0.1 \text{ GeV}$, we obtain for the two leptonic branching ratios,

$$\text{Br}(\tau \rightarrow e \bar{\nu}_e \nu_\tau a) = 1.2 \times 10^{-12} |g_{\tau\tau}|^2 \left(\frac{10^4 \text{ GeV}}{f_a} \right)^2, \quad (2.14)$$

$$\text{Br}(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau a) = 1.1 \times 10^{-12} |g_{\tau\tau}|^2 \left(\frac{10^4 \text{ GeV}}{f_a} \right)^2. \quad (2.15)$$

C. Other production modes

For universal LFC scenario, where $g_{\mu\mu} \sim g_{\tau\tau}$ one could consider including ALP production directly from D_s, D^0 and D^+ decays. For D_s the relevant channel is $D_s \rightarrow \mu \nu_\mu a$. As a numerical example let us take $m_a = 0.1 \text{ GeV}$, $f_a = 10^4 \text{ GeV}$, giving $\text{Br}(D_s \rightarrow \mu \nu_\mu a) \simeq 2.0 \times 10^{-13} |g_{\mu\mu}|^2$, which is comparable to $\text{Br}(D_s^+ \rightarrow \tau^+ \nu_\tau) \text{Br}(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau a) \simeq 2.1 \times 10^{-13} |g_{\tau\tau}|^2$ [using Eq. (2.15) and $\text{Br}(D_s^+ \rightarrow \tau^+ \nu_\tau) \simeq 5.4\%$].

For direct production of ALPs from $D^\pm$ and D^0 decays, the relevant channels are the exotic semileptonic decays $D \rightarrow K^{(*)} \mu \nu a$. For instance, $\text{Br}(D^+ \rightarrow \bar{K}^0 \mu^+ \nu_\mu a) = \mathcal{O}(10^{-14}) \times |g_{\mu\mu}|^2$ for $f_a = 10^4 \text{ GeV}$ and $m_a = 0.1 \text{ GeV}$. Since $\sim 10\%$ of the charm quarks fragment into D_s and the rest to $D^\pm$ and D^0 (sometimes via decays of excited charm mesons), the ALP production rate from these types of decays is again comparable to the one from τ decays.

As we will demonstrate later, an LFC ALP with $m_a > 2m_\mu$ is too short lived when $g_{\mu\mu} \sim g_{\tau\tau}$, making the limits and sensitivities discussed here largely irrelevant in this flavor-universal scenario. Below the dimuon threshold, the small-coupling parameter space is largely excluded anyway, so we neglect this contribution altogether, which is conservative. Note that we also ignore the direct production of ALPs from $D_s \rightarrow \tau \nu_\tau a$ decays since this channel is severely phase-space suppressed.

Since in the LFC scenarios the ALP production in τ decays is suppressed by both the Fermi constant and axion coupling, $\Gamma_{\tau \rightarrow \chi a} \propto \mathcal{O}(G_F^2 f_a^{-2})$, one may wonder if loop-induced processes may lead to enhanced production of ALPs in the FCNC decays of B and K mesons. If, for example, there is a direct coupling between an ALP and the W -boson pseudoscalar density $W_{\mu\nu}^a \tilde{W}^{a\mu\nu}$ [18], or if such a coupling is induced by the triangular lepton loop diagrams [19,20], there can be an enhanced production of ALPs due to one- or two-loop induced $(\partial_\mu a) \bar{d}_L \gamma_\mu s_L$ and $(\partial_\mu a) \bar{s}_L \gamma_\mu b_L$ operators. Such loop-suppressed production modes, e.g. $\Gamma_{K \rightarrow \pi a}$, can bypass the G_F^2 suppression and occur at the $\mathcal{O}(f_a^{-2})$ order. However, this type of production and ensuing constraints are model dependent—their evaluation requires a more detailed knowledge about the UV completion of the effective ALP interactions that we consider. While an effective $aW\tilde{W}$ coupling is inevitably induced due to triangular loop diagrams of τ and ν_τ , for example, it will have a rather strong form-factor suppression of $\mathcal{O}(m_\tau^2/m_W^2)$ inside the FCNC loops and these diagrams alone will not lead to additional constraints on the class of models considered here. Therefore, we concentrate on the tree-level production of ALPs, which are less susceptible to the unspecified model-building choices.

D. ALP decay

The dominant decay channels of the ALP depend heavily on its mass and the assumed flavor structures of the couplings. Figure 3 shows the dominant branching ratios of a leptophilic ALP as a function of m_a for the LFV anarchy (top), LFC flavor universal (middle), and LFC $e - \tau$ coupled (bottom) benchmark scenarios. We give the predictions for the individual partial decay widths below.

For an ALP decaying to a pair of leptons of the same flavor the decay width is given by

$$\Gamma(a \rightarrow \ell^+ \ell^-) = |g_{\ell\ell}|^2 \frac{m_\ell^2 m_a}{8\pi f_a^2} \left(1 - \frac{4m_\ell^2}{m_a^2}\right)^{1/2}, \quad (2.16)$$

while for decays to a pair of leptons of differing flavor the partial decay width is given by,

$$\begin{aligned} \Gamma(a \rightarrow \ell_1^\pm \ell_2^\mp) &= |g_{\ell_1 \ell_2}|^2 \frac{(m_{\ell_1} + m_{\ell_2})^2 m_a}{16\pi f_a^2} \\ &\times \left[1 - \frac{(m_{\ell_1} - m_{\ell_2})^2}{m_a^2}\right] v(m_a, m_{\ell_1}, m_{\ell_2}). \end{aligned} \quad (2.17)$$

Here, we define $a \rightarrow \ell_1 \ell_2$ as a sum of both modes, $a \rightarrow \ell_1^+ \ell_2^-$ and $a \rightarrow \ell_1^- \ell_2^+$, so that $\Gamma(a \rightarrow \ell_1^\pm \ell_2^\mp) = \Gamma(a \rightarrow \ell_1^+ \ell_2^-) + \Gamma(a \rightarrow \ell_1^- \ell_2^+)$. The ALP can also decay into two charged leptons and two neutrinos, e.g., $a \rightarrow \mu^+ \nu_\tau e^- \bar{\nu}_e$ through $g_{\mu\tau}$. However, these channels are suppressed by G_F^2 and by four-body decay phase space factors, and are irrelevant in our discussion.

The ALP-lepton couplings generate at one loop the $a \rightarrow \gamma\gamma$ decays. For our analysis these are phenomenologically important only in the corners of parameter space. Summing all lepton contributions for the $a \rightarrow \gamma\gamma$ decay width, we get

$$\Gamma(a \rightarrow \gamma\gamma) = \frac{\alpha^2 m_a^3}{64\pi^3 f_a^2} \left| \sum_{\ell=e,\mu,\tau} g_{\ell\ell} \left[1 - \frac{4m_\ell^2}{m_a^2} I_{\gamma\gamma}^2\left(\frac{m_a}{2m_\ell}\right)\right] \right|^2, \quad (2.18)$$

where

$$I_{\gamma\gamma}(z) = \begin{cases} \arcsin(z) & \text{for } z \leq 1, \\ \frac{\pi}{2} + \frac{i}{2} \log \left[\frac{z + \sqrt{z^2 - 1}}{z - \sqrt{z^2 - 1}} \right] & \text{for } z > 1. \end{cases} \quad (2.19)$$

For $m_a \gg 2m_\ell$, the term with $I_{\gamma\gamma}$ in Eq. (2.18) becomes much smaller than 1. For instance, if the axion is μ -phobic then in the range $2m_e < m_a < 2m_\tau$ the ratio of the diphoton width to the dielectron width grows with axion mass, $\Gamma(a \rightarrow \gamma\gamma)/\Gamma(a \rightarrow ee) \sim \alpha^2 m_a^2 / 8\pi^2 m_e^2 \sim 8 \times (m_a/m_\tau)^2$, where we have taken advantage of the numerical coincidence that $m_e/m_\tau \approx \alpha/8\pi$. Thus, there is a window around the τ mass, $m_\tau/3 \lesssim m_a \lesssim 2m_\tau$, where the photon decay mode is dominant. This can be seen in the bottom panel in Fig. 3. For $m_a \ll 2m_\ell$ the photon partial width scales with a high power of the axion mass

$$\Gamma(a \rightarrow \gamma\gamma) = \frac{\alpha^2 g_{\ell\ell}^2}{96^2 \pi^3} \frac{m_a^7}{m_\ell^4 f_a^2} + \mathcal{O}(m_a^9). \quad (2.20)$$

So for light τ -philic axions, where this is the only available decay channel, the axion will also be long-lived.

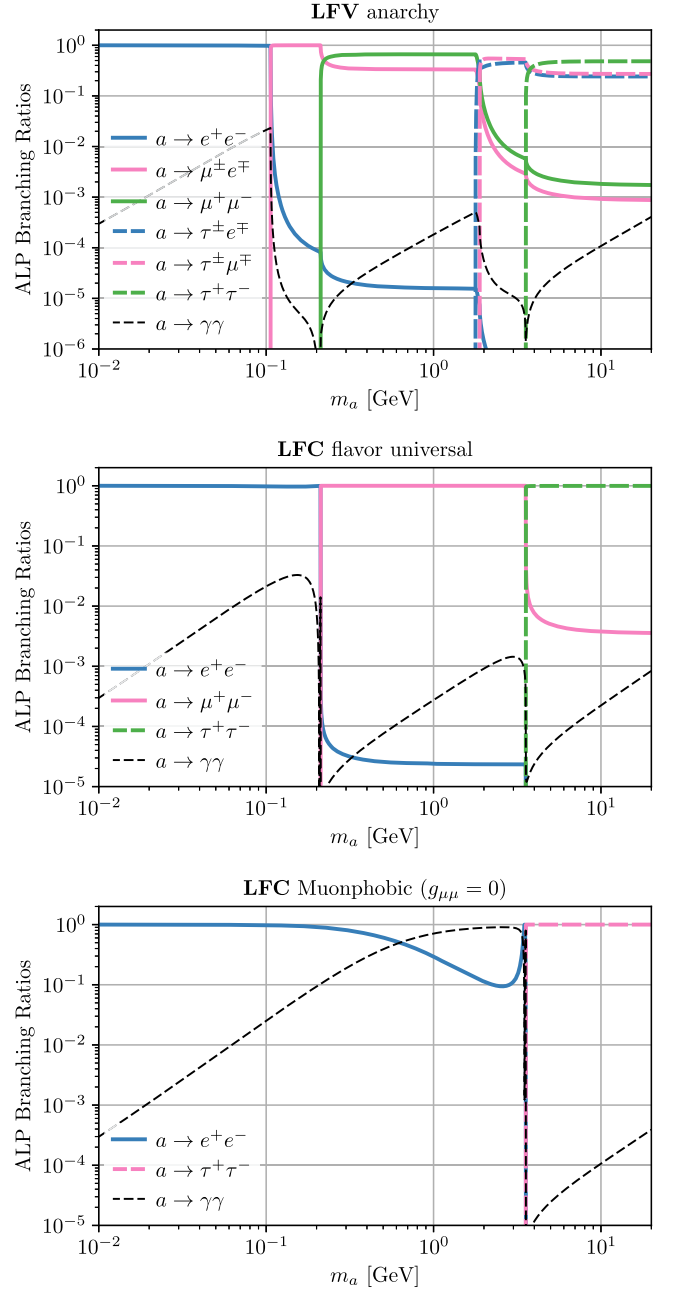


FIG. 3. The ALP decay branching ratios in the LFV anarchy (top), LFC flavor universal (middle), and LFC $e - \tau$ coupled (bottom) scenarios.

III. EXPERIMENTAL SIGNATURES

Next, we describe how we calculate the expected tau production rate and the resulting ALP signals at different experiments.

A. Tau production rate

We are particularly interested in the possible ALP fluxes from tau decays at three different proton beams:

- (1) NuMI at Fermilab: 120 GeV protons on a fixed target.
- (2) SPS at CERN: 400 GeV protons on a fixed target.
- (3) LHC: $\sqrt{s} = 13.6$ TeV center of mass (COM) energy of pp collisions.

Using PYTHIA8 [21] (v8.310 with the Monash tune choice of parameters [22]), we simulate individual pp collisions at each of the above three beam-target configurations. We checked that simulating pn collisions at fixed-target experiments gives similar τ yields and distributions.

We are interested primarily in the forward production of $D_{(s)}^\pm$, $\psi(2S)$, as well as direct $\tau^\pm$ production (the latter, though, is numerically subleading). Thus, we use the SoftQCD:all flag to simulate their production and turn off all HardQCD processes to avoid double counting. In each setup, we track the $D^\pm$ and $D_s^\pm$ mesons and $\psi(2S)$ charmonium produced by PYTHIA and force them to decay to $\tau^\pm$ leptons, weighting each event by the appropriate branching ratio to taus.

To speed up our simulations, we fit the transverse momentum p_T and Feynman x_F distributions of taus to the PYTHIA output using empirical analytical formulas. With the analytical fits, we can generate weighted tau events more efficiently, covering kinematical tails of distributions with smaller number of events and speeding up our parameter scan significantly. We do this for τ^+ and τ^- separately for all three beam setups above. Our fits also provide a simple description of tau production at the three beams that can be used in future applications.

In each detector, we load the same $\tau^\pm$ events generated with the analytical formulas and force taus to decay to ALPs for each of the new-physics decay channels discussed in Secs. II A and II B, keeping track of the weights and the corresponding branching ratios to ALPs.

In the estimates we neglect the interaction of secondaries within the target. While this component has been argued to

produce a non-negligible fraction of charmed mesons and tau particles in thick targets [23], it does not invalidate the bounds derived here. On the contrary, including this component can only boost the number of taus produced—our regions of coverage can be viewed as being conservative.

In our simulations we obtain the following inclusive $\tau^\pm$ yields per proton-nucleon (pN_T) collision,

- (1) NuMI: $2.9 \times 10^{-7} \tau^\pm / pN_T$.

- (2) SPS: $1.7 \times 10^{-6} \tau^\pm / pN_T$.

- (3) LHC: tau production cross section $\sigma_\tau^{\text{tot}} = 0.22$ mb.

Figure 4 displays the kinematic distributions of the produced $\tau^\pm$ leptons. The simulated samples correspond to approximately 2.2×10^6 τ 's for NuMI, 6.1×10^6 for SPS, and 2.1×10^6 for the LHC. In Appendix A, we discuss potential uncertainties in these predictions and their impact on the experimental sensitivities presented here. In Appendix B, we also provide a fit of our events with a parameterization for the differential cross section. This enables fast simulations of ALP production when larger simulation sample sizes are needed, but is not used in our main results.

We note that the issue of estimating the number of ALPs from tau decays at beam dump experiments is similar in nature to the calculation of the $\nu_\tau/\bar{\nu}_\tau$ flux. Both depend primarily on the production of $D_s^\pm$ mesons, and therefore, on an accurate description of open-charm production. At SHiP, there have been several estimates of this flux over the years [24–27], with total yields that vary by as much as an order of magnitude. In the future, the yields can be experimentally determined with significantly improved precision using data from the DsTau experiment (NA65 at the CERN SPS) [28], SHiP-charm [29], and with the neutrino data from SHiP itself [24], reducing the uncertainty on the ALP flux predictions. At FASER, the situation is similar and $\nu_\tau/\bar{\nu}_\tau$ data will help constrain it, as it already does for π , K production using $\nu_\mu/\bar{\nu}_\mu$ data [30]. The lower

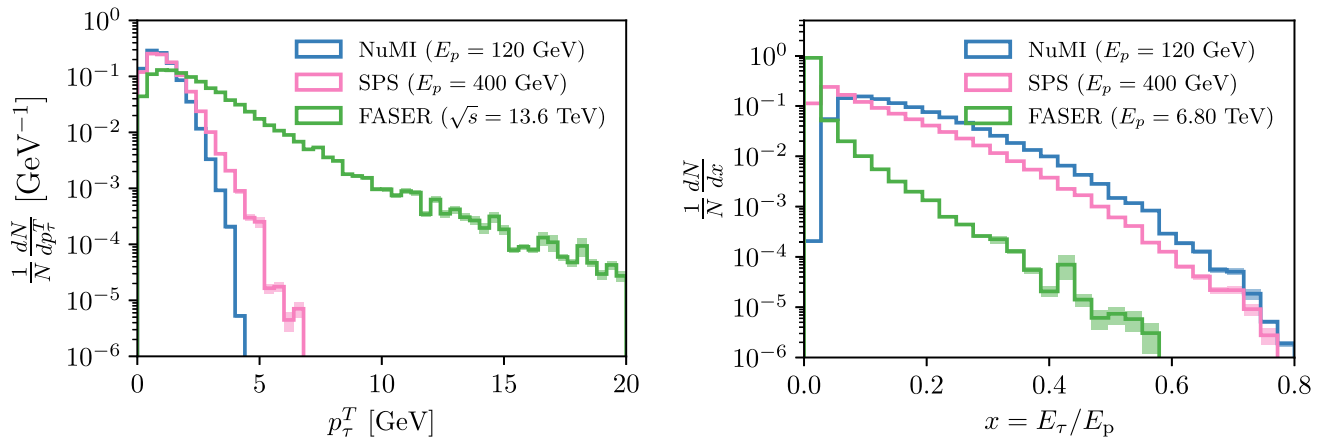


FIG. 4. Left: the normalized differential tau production rate at the collision point as function of the tau transverse momentum p_T^T (defined with respect to the z axis, i.e., the axis of the proton beam) for the three beams considered in this paper: NuMI, SPS, and FASER (as indicated). Right: the distribution of tau events as a function of the ratio between the total tau energy and the proton beam energy for the same three beams. The shaded bands show the statistical uncertainties on Monte Carlo samples.

COM energy at the NuMI beam means that there the uncertainty on the ALP flux is expected to remain large, given that the flux of $\nu_\tau/\bar{\nu}_\tau$ is too small to be observed. For more detailed studies on the impact of charm and tau production on new physics, see Refs. [31,32].

B. ALP flux and decays in flight

The ALP energy spectrum E_a in the tau rest frame is obtained by sampling the differential decay rates discussed in Sec. II. The ALP we consider here couples exclusively via the axial-vector current, so that its emission in the tau rest frame is isotropic, independent of the tau polarization [14]. In models where the ALP couples to left- or right-handed currents, then the decay is no longer isotropic for polarized taus. However, even in that case, the effect of the angular-dependence is small as the net polarization of the $\tau^+ + \tau^-$ produced at the target is small due to τ^+/τ^- production asymmetry being small. We find that the ratio of τ^+ to τ^- production rates at NuMI, SPS, and LHC is 0.73, 0.84, and 0.99, respectively. The dominance of D_s^- over D_s^+ at low energies is consistent with leading particle effects in charm hadroproduction [33].

The next step in our simulation is to calculate the expected number of decays in flight at different detectors. We assume that the detector's face is aligned with the line of sight to the target and then calculate the probability P_{decay} of an ALP to decay to a given final state $a \rightarrow \ell_1 \ell_2$ inside the fiducial volume as

$$P_{\text{decay}} = \mathcal{B}(a \rightarrow \ell_1 \ell_2) \exp \left[-\frac{\Gamma_a L}{\gamma_a \beta_a} \right] \left[1 - \exp \left[-\frac{\Gamma_a \lambda}{\gamma_a \beta_a} \right] \right], \quad (3.1)$$

where Γ_a is the total ALP decay width in the rest frame, γ_a and β_a are the boost factor and speed of the ALP, L is the distance from the target to the face of the detector, and λ is the distance that each ALP travels inside the fiducial volume. For long-lived ALPs ($\Gamma_a L \ll 1$ and $\Gamma_a \lambda \ll 1$), the expression in Eq. (3.1) simplifies to $P_{\text{decay}} = \lambda \Gamma_{a \rightarrow \ell_1 \ell_2} / \gamma_a \beta_a$, where $\Gamma_{a \rightarrow \ell_1 \ell_2}$ is the partial decay width for the $a \rightarrow \ell_1 \ell_2$ channel. For experiments whose decay volume is not instrumented (CHARM, SHiP, FASER, and FASER-2, see Table I), we furthermore estimate the probability for the decay products to reach the far side of the decay volume, after which these are assumed to be detected in the calorimeter. In all cases we apply simple energetic requirements on the final state leptons/photons. Beyond these geometric requirements—that axions enter the front of the fiducial volume, decay in the detector volume, and the leptons hit the electromagnetic calorimeter (for the experiments that require it)—and simple energetic cuts, we do not apply detector efficiencies.

C. ALP detection

The list of detectors we consider at the three locations is shown in Table I. For each, we quote the assumed fiducial volume, the position of the center of the detector face, and

TABLE I. List of experiments considered in this work and their respective parameters. Here, $\mathcal{L}$ stands for the assumed LHC luminosity, and $\delta x(\delta y)$ are the assumed locations of the centers of the detectors in the horizontal (vertical) directions. The dimensions of the detector are given as width $\times$ height $\times$ depth and the baseline is the distance from the point of τ production to the front of the detector. Experiments for which the axion decay products are required to hit an electromagnetic calorimeter at the back of the fiducial volume are denoted with “Ecal.”

Experiment	p^+ energy	Exposure	Baseline	Fiducial volume	Material(s)	Cuts	References
SHiP	400 GeV (SPS)	6×10^{20} POT	33.7 m	Pyramidal frustum ($1 \times 2.7 \rightarrow 4 \times 6$) $\times$ 49.6 m ³	He, Ecal	$E_{\ell, \gamma} > 0.5$ GeV	[34]
DUNE ND	120 GeV (NuMI/LBNF)	11×10^{21} POT (10 yrs)	574 m	($5 \times 5 \times 5.88$) m ³	LAr	$E_{\ell, \gamma} > 0.5$ GeV	[35]
protoDUNE NP02 [NP04]	400 GeV (SPS)	1.75×10^{19} POT (5 yrs)	677 [723] m T2	($6 \times 7 \times 6$) m ³ $\delta x = +3[-3]$ m, $\delta y = +1.5$ m	LAr	$E_{\ell, \gamma} > 0.5$ GeV	[36]
FASER	$\sqrt{s} = 13.6$ TeV (LHC)	$\mathcal{L} = 150$ fb ⁻¹	480 m	$R = 0.1$ m, $z = 1.5$ m $\delta y = 6.5$ cm	Air, Ecal	$E_{\ell, \gamma} > 50$ GeV	[37]
FASER2	$\sqrt{s} = 13.6$ TeV (LHC)	$\mathcal{L} = 3$ ab ⁻¹	650 m	$2.6 \times 1 \times 10$ m ³	Air, Ecal	$E_{\ell, \gamma} > 50$ GeV	[38]
ICARUS	120 GeV (NuMI)	3×10^{21} POT	803 m (5.6°)	$2 \times (2.67 \times 2.86) \times 17.0$ m ³	LAr	$E_{\ell, \gamma} > 0.1$ GeV	[39,40]
NOvA ND	120 GeV (NuMI)	4.2×10^{21} POT (2024)	990 m (14.6 mrad)	$3.9 \times 3.9 \times 12.7$ m ³	Liquid Scint.	$E_{\ell, \gamma} > 0.1$ GeV	[41]
MicroBooNE	120 GeV (NuMI)	2×10^{21} POT	685 m (8°)	$2.26 \times 2.03 \times 9.42$ m ³	LAr	$E_{\ell, \gamma} > 0.1$ GeV	[42]
BEBC	400 GeV (SPS)	2.72×10^{18} POT	404 m	$3.57 \times 2.52 \times 1.85$ m ³	Ne/H ₂	$ \vec{p}_{\ell, \gamma} > 0.5$ GeV	[43,44]
CHARM	400 GeV (SPS)	2.4×10^{18} POT	480 m	$3 \times 3 \times 35$ m ³ $\delta x = 5$ m	Air, Ecal	$E_{\ell, \gamma} > 0.5$ GeV	[45,46]

their known or expected exposure in terms of protons on target (POT) or luminosity. For FASER and FASER-2, we use the results from [47]. While a limit on ALP couplings from FASER data could in principle already be placed by recasting their latest search for dark photons [37], we find that this is not yet competitive with other direct limits. For ProtoDUNE, we follow closely Ref. [36]. The faces of the two detectors, NP02 and NP04, are not aligned with the direction of the beam. We expect that accounting for this misalignment would not significantly affect our results. We take into account τ production in the primary target and neglect the one from the beam dump, where a smaller, less-energetic fraction of the beam stops. For ICARUS and MicroBooNE, we follow Refs. [39,40], assuming a total NuMI POT of 2×10^{21} at MicroBooNE [48] and a projected 3×10^{21} POTs at ICARUS. We note that the polarity of the magnetic horns is not important for the ALP signal since it is produced by the promptly decaying $D_{(s)}$ and τ particles, and so we sum the so-called forward and reverse horn current data taking periods.

For the DUNE near detector (ND), one may consider a few different possibilities: ND LAr, ND GAr, or SAND, each with their unique set of advantages. For charged final states, ND GAr will likely be preferable from the point of view of backgrounds and efficiencies. For two-photon decays, however, a higher density detector would be needed, so ND LAr or SAND would be more relevant. In [49], the authors find large detection efficiencies for ND GAr, above 90% for the GAr module. For SHiP, we take the numbers from [34], quoting the sensitivity for the benchmark of 6×10^{20} POTs.

Our results are quoted in terms of event-rate sensitivities, i.e., the regions where experiments can expect more than 2.3 signal events during their total exposure. In a background-free search, this would correspond to the 90% CL sensitivity. The zero (or $\ll 1$) background limit is applicable to the existing searches at CHARM and BEBC. In these two cases one can therefore interpret our results as the 90% CL exclusions. For future searches, whether or not the zero background limit will apply will depend on the details of the experiments and the analysis strategies adopted. For instance, in the existing searches for dilepton pairs at MicroBooNE [42,50,51] and ICARUS [40], requiring the dilepton pair invariant mass to coincide with the assumed axion mass reduces backgrounds to be below or at the level of one event, with efficiencies that range from 10%–50%. The ALP signals discussed here are more energetic than, for instance, the signals one would obtain from LLPs that would be produced in decays of kaons at rest. As such, we expect the neutrino-induced backgrounds to be less important in the case of ALPs from tau decays.

On the other hand, note, however, that the reconstruction of dimuon pairs may worsen compared to LLPs, since the muons are now more likely to escape the fiducial volume,

although it opens up the possibility to search for nonstarting ALP decay events. Understanding how big this effect is, however, requires a rather detailed analysis of the backgrounds and the selections, which we expect to be done by the collaborations as part of their search strategy optimization.

D. Other experiments

Let us also comment on several potentially relevant experiments that we do not include in the figures. The ArgoNeuT detector at the Fermilab MINOS near detector hall has searched for heavy neutral leptons and axionlike particles using dimuon pairs starting outside or inside of its Liquid Argon Time Projection Chamber (LArTPC) fiducial volume [52,53]. In Ref. [54], the authors adapted these searches to place new limits on the parameter space of leptophilic ALPs and Majorons. We take 87% of ALPs to be produced at the NuMI target and 13% at the location of the beam dump, counting as signal all ALP decays to dimuons that take place within a $40 \times 47 \times 210 \text{ cm}^3$ volume that starts 1033 m away from the target and is centered around the proton beam. With a simplified analysis which takes the effect of geometric cuts and signal efficiency to be independent of energy and tuning their overall effect ϵ , we find some agreement with Bertuzzo *et al.* [54] when $\epsilon \approx 0.5$ but were unable to fully reproduce their results. We find some agreement with the limits on $1/f_a$ from Ref. [54] for the case $R_a = g_{\ell\ell}/g_{\ell\ell_2} = 5$, but not for $R_a = 1/3$. Nevertheless, we find that ArgoNeuT is not competitive with existing limits placed by CHARM and BEBC searches, even when using the exclusion curves in Ref. [54], or when using the most optimistic assumptions about the efficiencies and exposure.

Recently, the 2×2 DUNE near detector demonstrator, also a LArTPC detector, has been operating in the MINOS hall [55,56]. It consists of four tall detectors, each $0.7 \times 1.4 \times 0.7 \text{ m}^3$ of volume and located at a similar distance to the NuMI target as ArgoNeuT. We find that for 100% efficiencies and no backgrounds, a search for starting events (decays inside the LArTPCs) would not be competitive with existing limits or with the event-rate sensitivities of MicroBooNE, ICARUS, and NOvA even for an exposure of 10^{21} POTs.

The NA62 kaon detector has also been used to set limits on long-lived particles using the SPS beam [57,58]. We estimate that in the future, with 10^{18} POT [59], NA62 in beam dump mode may outperform existing CHARM and BEBC limits on the ALP models we are interested in. Finally, the past NuCal detector at the 70 GeV proton beam at IHEP could also be potentially relevant [60,61]. While we do not study its reach in detail, we observe that the smaller $\tau^\pm$ yield due to the lower energy of the beam could be compensated by the closer location of the detector (63 m) and its relatively large fiducial volume ($2.4 \times 2.4 \times 23 \text{ m}^3$) and 1.7×10^{18} POTs (see, for instance, Refs. [62,63]).

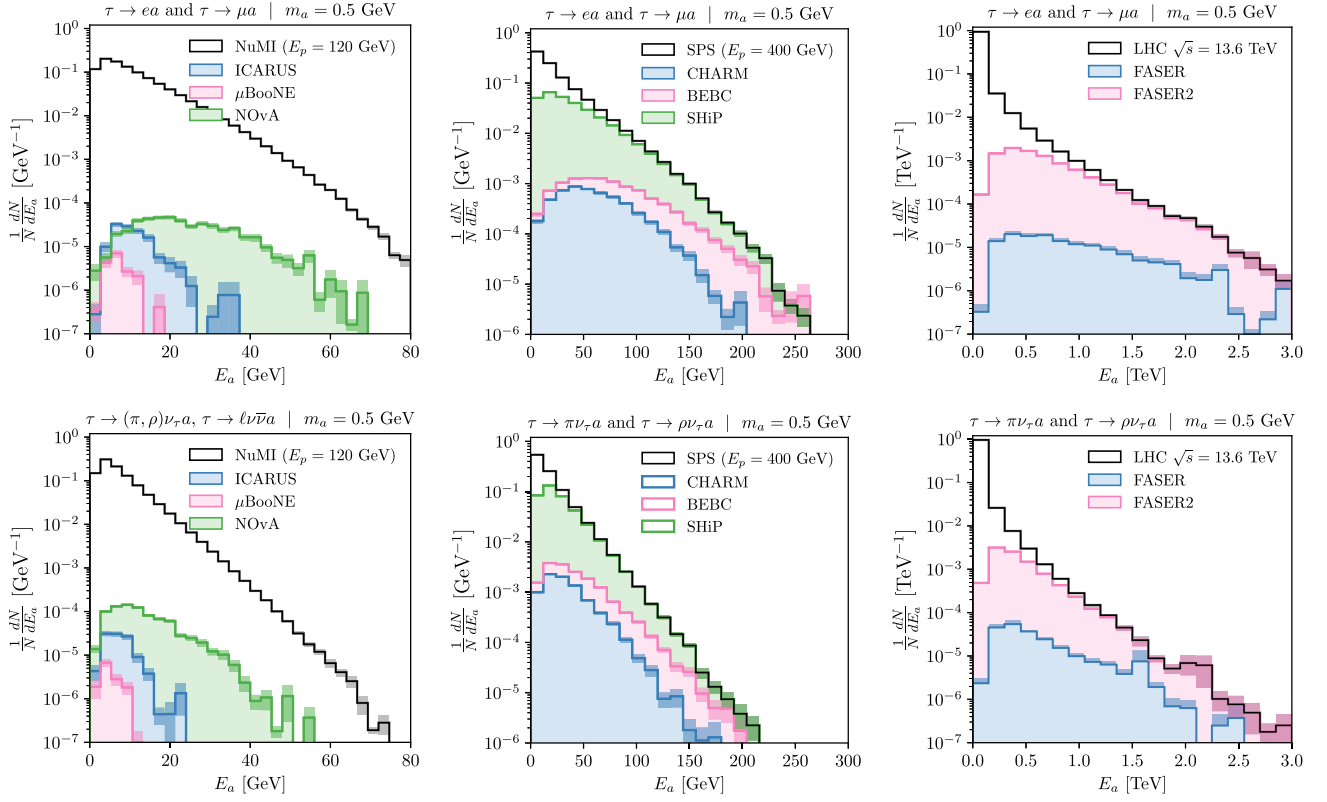


FIG. 5. Energy distribution of ALPs produced in LFV (top row) and LFC (bottom row) decays of taus at NuMI (left column, 120 GeV proton beam), SPS (center column, 400 GeV proton beam), and at the LHC (right column, $\sqrt{s} = 13.6$ TeV pp collisions). The solid black lines show the total spectrum at the production point and the filled histograms show the respective fraction of the spectrum within the acceptance of each experiment. For all panels, we fix $m_a = 0.5$ GeV. The uncertainty from the Monte-Carlo statistics is shown as shaded envelopes.

IV. RESULTS

Next, we present our results for the experimental reach for each of the benchmark LFV and LFC ALP scenarios introduced in Sec. I, see Eqs. (1.7) and (1.12).²

A. Flavor-violating ALPs

For LFV ALPs, we consider two benchmarks, the anarchic flavor pattern, Eq. (1.7), and the flavor hierarchical case, Eq. (1.8). The exclusions and the projected event-rate sensitivities for the flavor anarchy LFV case, Eq. (1.7), are shown in Fig. 6. The left panel shows as filled colored regions in the (m_a, f_a^{-1}) plane the bounds we derived from CHARM and BEBC searches. These are compared to other direct limits on ALPs from muon and tau decays as well as supernovae, shown in shades of grey (see Sec. IV C). The projected event-rate sensitivities of various experiments are

²In the LFV anarchy and LFC universal cases, we have checked that our results are in reasonable agreement with SensCalc [63,64] when using the same τ yields obtained in our simulation. We note that SensCalc includes the secondary charm production at SPS and therefore uses different τ distributions, and approximates the four-body τ decay matrix element to be constant.

shown as lines of different colors and styles, as indicated in the legend.

The right panel in Fig. 6 shows an enlarged region of the parameter space. Despite the lower energy of the main injector beam compared to SPS and the LHC, neutrino detectors off axis from the NuMI target, namely NOvA, MicroBooNE, and ICARUS, may improve on existing limits, if a high-efficiency, low-background search can be performed. The SBND detector could, in principle, also be considered, however, it is much farther off axis than the other three. We note that in the models we are interested in the ALP signatures are typically of relatively high energy (see Fig. 5), so that a search for a nonstarting dimuon signal could be advantageous, further increasing the effective fiducial volume of the experiment, as already demonstrated by the ArgoNeuT experiment in Refs. [52,53].

Figure 7 shows our results for the flavor hierarchical cases, Eq. (1.8). Each panel corresponds to a different hierarchy between LFV and LFC couplings; the flavor diagonal couplings are set to $g_{\ell\ell} = 1$, while the values of off-diagonal couplings $g_{\mu\tau}$, $g_{e\tau}$, and $g_{e\mu}$ decrease from top left to bottom right panel. For simplicity, all off-diagonal couplings were set to be equal, $g_{\mu\tau} = g_{e\tau} = g_{e\mu}$, in the four

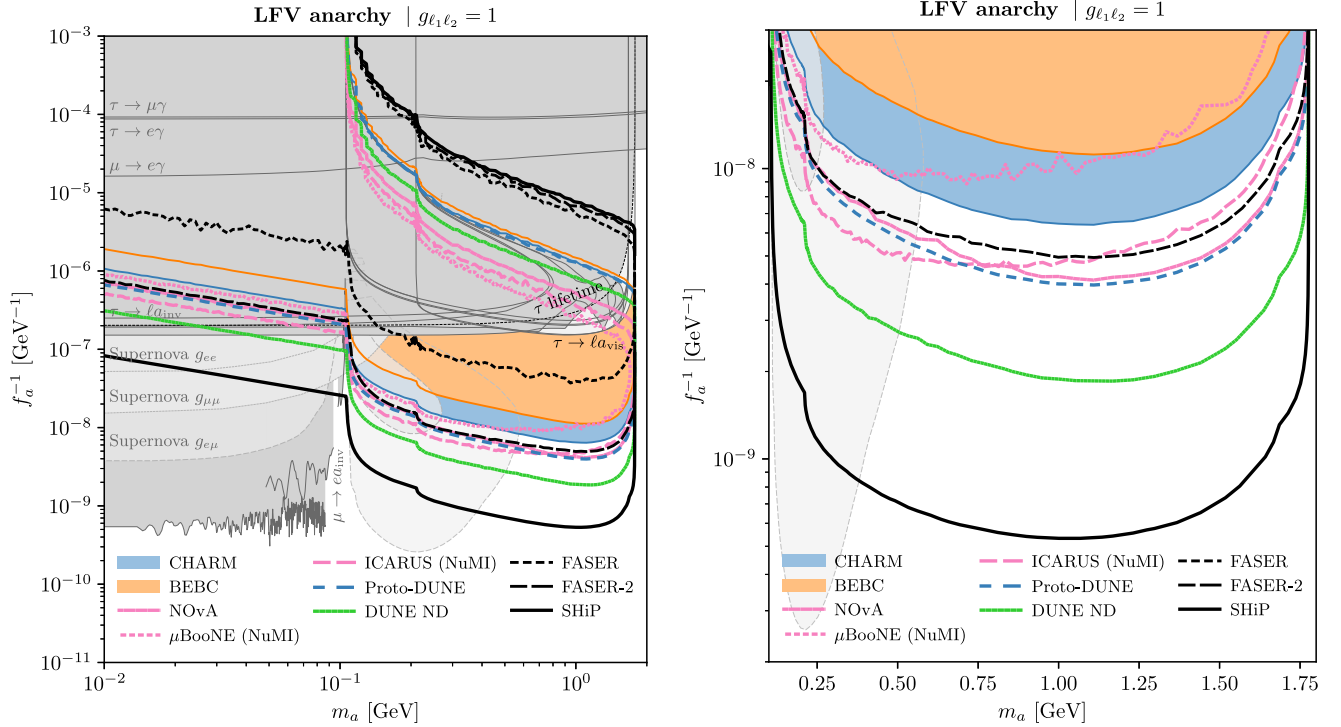


FIG. 6. Parameter space of flavor-violating ALPs with flavor anarchical couplings, Eq. (1.7). The right panel shows an enlarged version of the parameter space shown in the left panel. Constraints on long-lived ALPs discussed in Sec. III are shown as filled regions for existing limits from CHARM and BEBC and as solid and dashed lines for projections for neutrino experiments, FASER, proto-DUNE, and SHiP. Other constraints from tau decays, discussed in Sec. IV, are shown as light grey regions for $\tau \rightarrow \ell a_{\text{inv}}$ and dark grey regions for $\tau \rightarrow \ell a_{\text{vis}}$.

representative cases shown. That is, the variation of LFV couplings is modeled using a single parameter $\lambda \ll 1$, setting,

$$\begin{pmatrix} g_{ee} & g_{e\mu} & g_{e\tau} \\ g_{e\mu} & g_{\mu\mu} & g_{\mu\tau} \\ g_{e\tau} & g_{\mu\tau} & g_{\tau\tau} \end{pmatrix} = \begin{pmatrix} 1 & \lambda & \lambda \\ \lambda & 1 & \lambda \\ \lambda & \lambda & 1 \end{pmatrix}. \quad (4.1)$$

The decreasing values of off-diagonal couplings in the panels in Fig. 7 illustrate the effect of a decreasing amount of LFV on the ALP phenomenology. In practice, this variation amounts to changing the ALP production rates without substantially modifying the ALP lifetime, the latter being driven by $g_{ee,\mu\mu}$. Consequently, the experiments that are sensitive to shorter-lived ALPs, such as FASER and SHiP, become increasingly more relevant.

From low energy perspective, the exact choice for the values of the off-diagonal ALP couplings in Eq. (4.1) is arbitrary, since it depends on the origin of flavor structure and thus on details of UV physics. Changing the relative sizes of $g_{e\mu}$ and $g_{e(\mu)\tau}$ couplings does not change appreciably the projected sensitivities of beam-dump-like experiments. It does, however, directly control the importance of charged lepton flavor violation searches. For the choice of equal size LFV couplings, $g_{e\mu} = g_{e(\mu)\tau}$, the constraints from $\mu \rightarrow ea$ decays are typically the dominant laboratory

constraints whenever $m_a < m_\mu - m_e$. This need not be the case if, for instance, we had assumed $g_{e\mu} \ll g_{e(\mu)\tau}$. Nevertheless, for the LFV hierarchical case, the choice of $g_{e\mu}$ has little impact on ALP searches at high-energy beam dumps in the most relevant mass region, namely $m_a > 2m_\mu$. In the intermediate region of $m_e + m_\mu < m_a < 2m_\mu$, the LFV ALP decay channel can be important depending on the size of $g_{\mu e}/g_{ee}$.

B. Flavor-conserving ALPs

The left panel of Fig. 8 shows our results for the flavor-conserving scenario, Eq. (1.9), with universal couplings to all three generations of SM leptons, $g_{ee} = g_{\mu\mu} = g_{\tau\tau}$. For ALP production due to ALP bremsstrahlung modifying the SM τ decays, the additional G_F suppression of the decay amplitude compared to the LFV case means that the sensitivity to scale f is significantly weaker. Existing limits from CHARM and BEBC based on ALP production in tau decays are comparable with the other beam dump limits shown in various shades of grey. The latter are based on electron bremsstrahlung and light meson decays. Given the larger number of electrons and mesons available at the beam dump facilities, these searches can still be relevant to the LFC scenarios, despite the strong suppression of ALP couplings to electrons. We find that the other considered

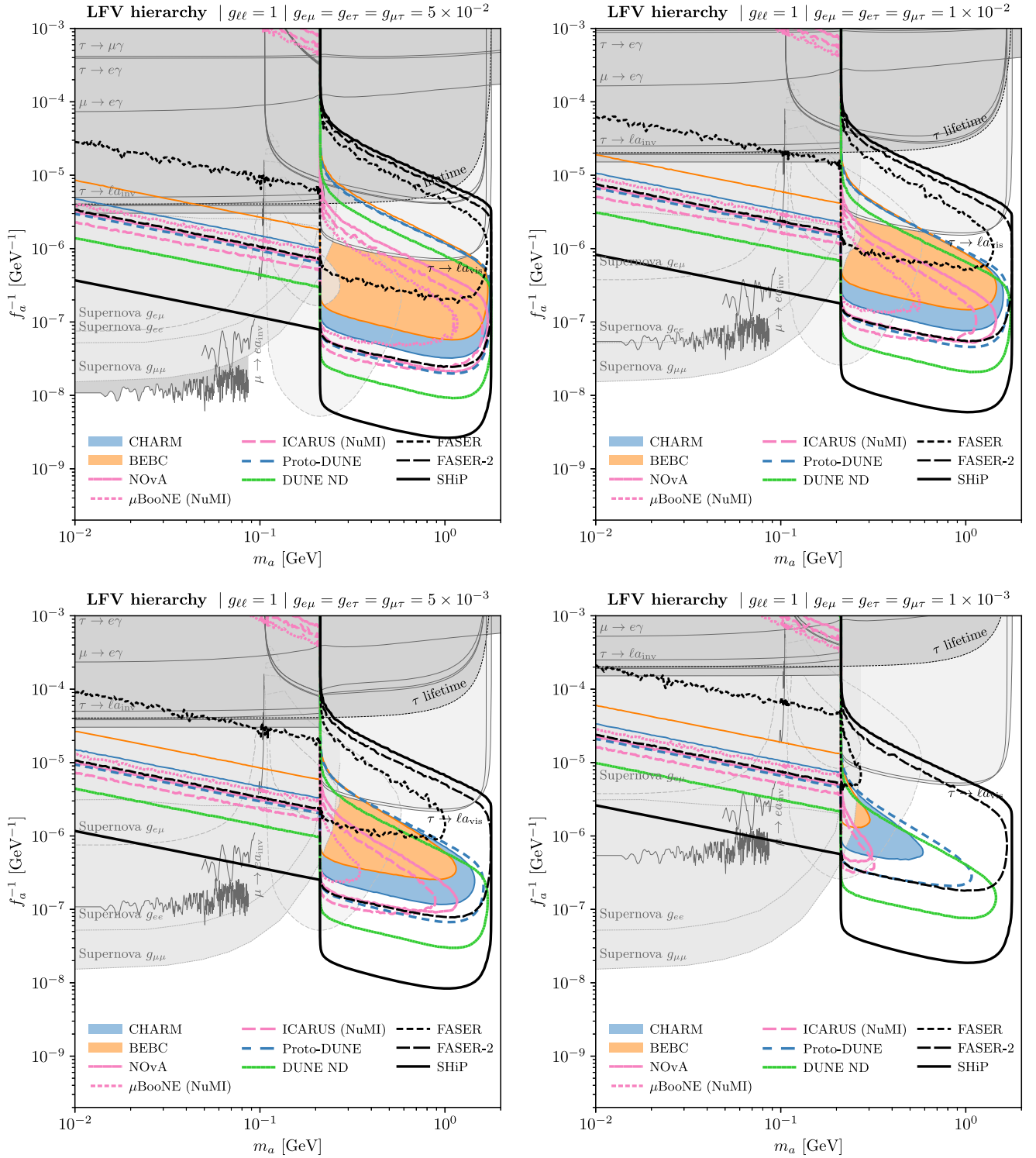


FIG. 7. Same as Fig. 6, but for flavor-violating hierarchical couplings, Eq. (1.8), for several representative sizes of flavor off-diagonal couplings, as indicated. Constraints on long-lived ALPs discussed in Sec. III are shown as filled colored regions for existing limits from CHARM and BEBC and as solid and dashed lines for projections for neutrino experiments, FASER, proto-DUNE, and SHiP. Other constraints from tau decays, discussed in Sec. IV, are shown as light grey regions for $\tau \rightarrow \ell a_{\text{inv}}$ and dark grey regions for $\tau \rightarrow \ell a_{\text{vis}}$.

experiments, MicroBooNE, ICARUS, NOvA, and FASER, are not competitive with the electron-based limits. In the future, however, the high-energy experiments such as

FASER-2 and SHiP will be able to explore new parameter space using only ALP production from the tau decays we calculate here.

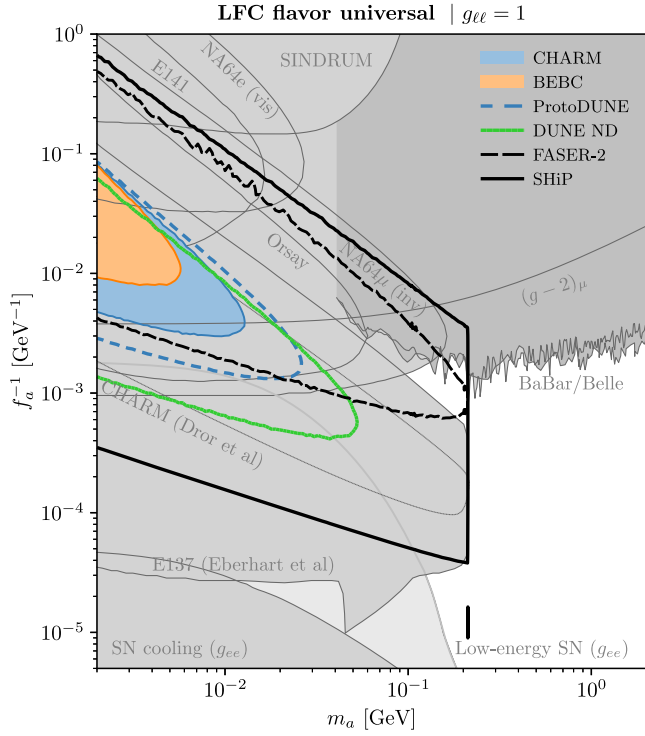


FIG. 8. Limits and projected event-rate sensitivities in the parameter space of LFC ALP with universal couplings of Eq. (1.9). Constraints on long-lived ALPs discussed in Sec. III are shown as filled regions for existing limits from CHARM and BEBC and as solid and dashed lines for projections for FASER-2, proto-DUNE, DUNE, and SHiP. Other neutrino experiments and FASER are less sensitive and not shown. The grey regions show limits from SINDRUM [67] and CHARM [45] as derived in [15], electron beam dumps limits derived in as derived in [68,69] (see also Refs. [70,71]) with data from E137 [72], E141 [73], and Orsay [74], as well as limits from BABAR [75], Belle [76], and NA64 [77]. We also show supernova limits based on cooling arguments [78,79] as well as the constraint from ALP energy depositions in the outer envelope of the supernova [79,80] as a lighter grey line.

Figure 9 shows the results for the two flavor nonuniversal LFC benchmarks, $g_{ee} = g_{\tau\tau} = 1, g_{\mu\mu} = 0$ (μ -phobic, left panel), and $g_{ee} = g_{\mu\mu} = 0, g_{\tau\tau} = 1$ (τ -philic, right panel). In both benchmarks the ALP does not couple to muons, which has important phenomenological consequences. For instance, the flavor universal LFC ALP, Fig. 8, with a mass above the dimuon threshold, $m_a > 2m_\mu$, would decay via $a \rightarrow \mu^+\mu^-$ with a lifetime too short to be observed in beam dump experiments. For vanishing coupling to muons, $g_{\mu\mu} = 0$, the ALP is long-lived also for masses above the dimuon threshold, decaying either via $a \rightarrow e^+e^-$ or $a \rightarrow \gamma\gamma$, and thus beam dump searches are sensitive to the two scenarios also above the muon threshold. In fact, for a τ -philic ALP with $g_{ee} = g_{\mu\mu} = 0$, only $a \rightarrow \gamma\gamma$ decays are possible below the di-tau threshold, $m_a < 2m_\tau$, making the τ -philic benchmark significantly less constrained than the

μ -phobic one. Note also, that in the τ -philic case there is a steeper dependence of the sensitivities to ALP mass because the photon width scales with a high power of ALP mass, $\Gamma(a \rightarrow \gamma\gamma) \sim m_a^7$, see Eq. (2.20). In summary, as shown in Fig. 9, ALP production from tau decays at beam dumps are again providing some of the leading constraints in the parameter space also for the flavor nonuniversal LFC ALPs.

C. Existing limits

In Figs. 6–9 we compare the reach of beam dump experiments with other probes of weakly interacting particles. Because the ALP couples exclusively to leptons, many limits on long-lived ALPs from pion or kaon decays, such as [53,65,66], are too weak to appear in our plots. This is particularly true for LFV couplings, where the ALP production in fully leptonic two-body τ and μ decays is significantly enhanced relative to any G_F -suppressed production mode from pion and/or kaon decays. The tau lifetime measurements, bounds on some of the rare tau decay channels, and the astrophysical constraints therefore pose more meaningful limits on the ALP decay constant. For the LFC case, the tau and muon ALP-branching ratios are suppressed by G_F and are thus much smaller than in the LFV case for the same value of f_a . In this case direct constraints, therefore, come mostly from ALP couplings to electrons.

Two-body decays of taus provide some of the leading direct constraints on the leptophilic ALPs. Decays into invisible bosons, $\tau \rightarrow \ell a_{\text{inv}}$, have been carried out by ARGUS [87], Belle [88], and Belle-II [89] and constrain these BR's to be as low as $\mathcal{O}(10^{-3})$. When the ALP is short-lived, it leads to tau decays to neutrinoless three-lepton final states, such as $\tau^+ \rightarrow \mu^+\mu^+\mu^-$. We refer to this case as $\tau \rightarrow \ell a_{\text{vis}}$ bounds. These are constrained by BABAR [90], Belle [91], ATLAS [92], CMS [93], LHCb [94] and Belle-II [95]. In setting the constraints we use the world best limit as collected by PDG [10],

$$\mathcal{B}(\tau^- \rightarrow e^- e^+ e^-) < 2.7 \times 10^{-8}, \quad (4.2)$$

$$\mathcal{B}(\tau^- \rightarrow \mu^- \mu^+ \mu^-) < 1.9 \times 10^{-8}, \quad (4.3)$$

$$\mathcal{B}(\tau^- \rightarrow \mu^- e^+ \mu^-) < 1.7 \times 10^{-8}, \quad (4.4)$$

$$\mathcal{B}(\tau^- \rightarrow e^- \mu^+ \mu^-) < 2.7 \times 10^{-8}, \quad (4.5)$$

$$\mathcal{B}(\tau^- \rightarrow e^- e^+ \mu^-) < 1.8 \times 10^{-8}, \quad (4.6)$$

$$\mathcal{B}(\tau^- \rightarrow e^- \mu^+ e^-) < 1.5 \times 10^{-8}. \quad (4.7)$$

These are dominated by constraints from BABAR, Belle, and Belle-II, where tau production proceeds via $e^+e^- \rightarrow \tau^+\tau^-$. When setting constraints on leptophilic ALP parameter space we therefore require that the ALP decays within 1 cm, assuming an average energy of

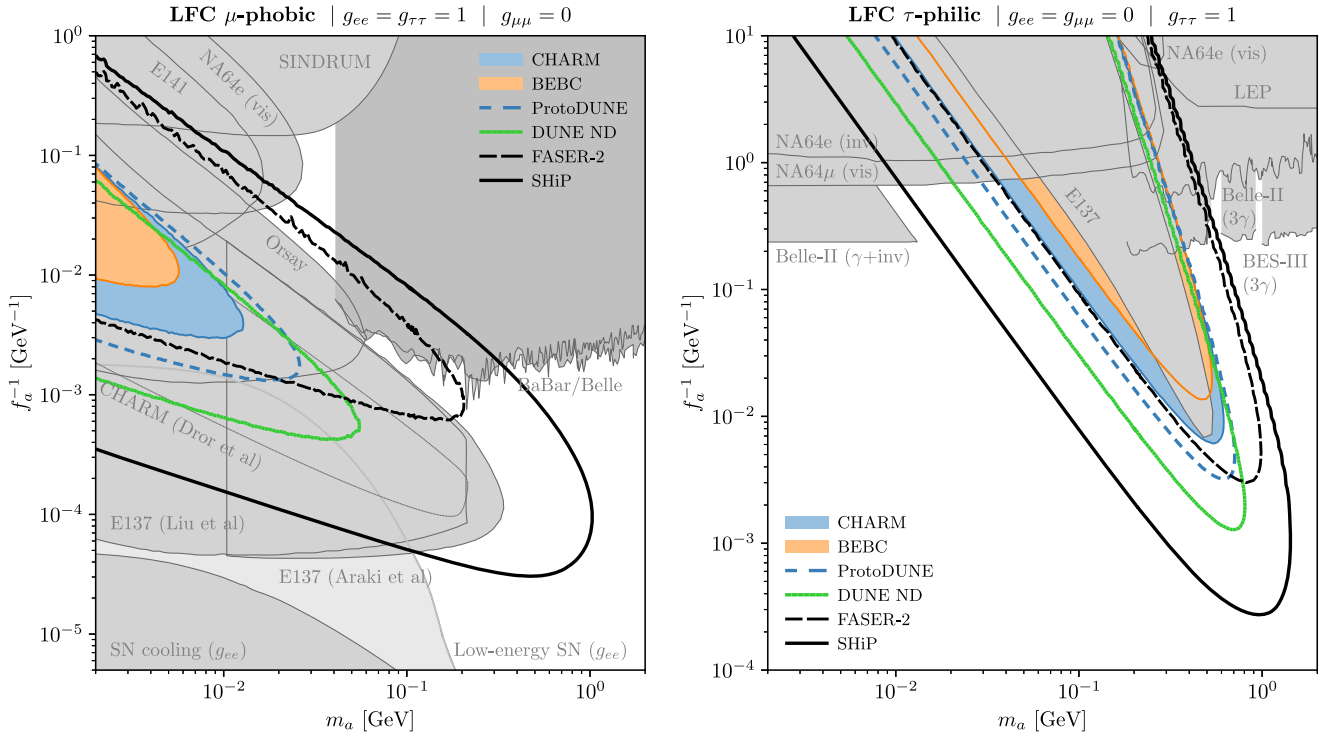


FIG. 9. Same as Fig. 8 but for two flavor nonuniversal benchmarks a μ -phobic ALP (left) and a τ -philic ALP (right), cf. Eqs. (1.10) and (1.12). The grey regions in the right panel show the limits derived in [81] using data from E137 [72], NA64 [82–84], and LEP [85] as well as limits from Belle-II and BES-III derived in [86].

$$\langle E_a \rangle \sim \frac{\sqrt{s}}{4} \left(1 + \frac{m_a^2}{m_\tau^2} - \frac{m_\ell^2}{m_\tau^2} \right), \quad (4.8)$$

where m_ℓ is the mass of the lepton produced in $\tau \rightarrow \ell a$ decays. For $\tau \rightarrow \ell a_{\text{inv}}$, we require that the ALP does not decay within 3 m. The excluded regions are shown in light (dark) grey in Figs. 6 and 7 for visible (invisible) ALPs.

The tau lifetime is well known in the SM. The leptonic and hadronic branching ratios of the τ are known, in some cases, to N³LO [96–98], while the experimental determinations have $\sim 0.1\%$ precision [10]. The fit of all observed branching ratios is consistent with the SM at the per-mille level [10]. We use this to impose a constraint of $\mathcal{B}(\tau \rightarrow \ell a) < 10^{-3}$, which is independent of the decay properties of a . This constraint is shown as a dashed black line in Figs. 6 and 7.

Flavor-violating muon decays can also be directly constrained. Experimental bounds on $g_{e\mu}/f_a$ are derived from $\mu \rightarrow ea$ decays, with ALP being either visible or invisible. In our plots, $\mu \rightarrow ea_{\text{inv}}$ bounds are the most relevant, since for such small masses and weak couplings the ALP is almost always long-lived for $m_a < m_\mu - m_e$. We show the experimental limits from TWIST [12], PIENU [13], and Refs. [99,100] (limits from Jodidio *et al.* [11] are only valid below 12 MeV). These curves are the analogs of the isotropic case shown in Ref. [14].

Off-shell ALP exchange in $\mu \rightarrow e\gamma$, $\mu^+ \rightarrow e^+e^+e^-$, and $\mu \rightarrow e$ conversion, as well as analogous channels for taus, also constrain leptophilic ALPs. In the region of interest for this study, these constraints are subdominant relative to direct emission in lepton decays, however, they can dominate for $m_a > m_\tau - m_e$. For comparison, in our figures we show the limits on f_a from lepton decays to single photons such as $\mu \rightarrow e\gamma$, following Ref. [101]. Recently, MEG-II [102] placed a world-leading limit on $\mathcal{B}(\mu^+ \rightarrow e^+\gamma) < 1.5 \times 10^{-13}$ at 90% CL, constraining decay constants as large as $f_a \gtrsim \mathcal{O}(10^5)$ GeV in a LFV anarchical scenario. This is to be compared with lower limits of $f_a \gtrsim \mathcal{O}(10^7)$ from tau decays such as $\tau \rightarrow ea$. Tau decays to a lepton and a photon benefit from the larger ALP couplings, but, for flavor-anarchical scenarios, still provide weaker constraints due to the weaker limits on the branching ratios. Here, we show limits from BABAR, $\mathcal{B}(\tau \rightarrow e\gamma) < 3.3 \times 10^{-8}$ [103], and Belle, $\mathcal{B}(\tau \rightarrow \mu\gamma) < 4.2 \times 10^{-8}$ [104], both at 90% CL.

Astrophysical constraints are relevant at larger decay constants (lower couplings) and are most effective for ALP masses below $m_a \sim \mathcal{O}(100)$ MeV. In particular, the LFC couplings to electrons and muons are strongly constrained by the observation of SN1987A (see, e.g., Refs. [78–80]). There are multiple mechanisms to derive a limit, depending on the imposed physical requirement. The most well-known is the cooling argument, based on the fact that weakly coupled ALPs can be produced and escape the star,

carrying away energy from deep inside the progenitor. Given that the observed neutrinos from SN1987A are consistent with the energy loss expected from SM processes within a protoneutron star, an additional dominant cooling mechanism is ruled out. We show limits on g_{ee}/f_a and $g_{\mu\mu}/f_a$ based on this criterion directly from [14] as a lighter shade of grey and in dashed lines. Supernova cooling limits on the LFV coupling $g_{e\mu}$ from [105] are also shown in the same style. The $g_{e\mu}$ coupling allows for $e - \mu$ coalescence production of ALPs in the protoneutron star, which is the dominant process for SN cooling constraint on this coupling for ALP masses in the range $115 \text{ MeV} \lesssim m_a \lesssim 280 \text{ MeV}$. Note that in the figures we plot each of these bounds (labeled as “supernova $g_{ee,\mu\mu,e\mu}$ ”) individually, where we consider each of the $g_{\ell\ell'}$ couplings is nonzero at a time. This is not entirely correct, for each flavor pattern one should recompute the SN cooling bounds taking all the processes into account simultaneously. However, since one of the couplings typically dominates, this suffices for our purposes—indicating that there is significant open parameter space for leptophilic ALPs.

Cooling is not the only way to constrain ALPs with supernovae, however. For instance, Ref. [106] derives limits on g_{ee}/f_a by arguing that the energy outflow induced by the propagating but “trapped” ALPs would stall the supernova explosion. These constraints are sensitive to the exact details of the supernova model, and can thus be viewed as less conservative (i.e., the exact shape of the constraints is more dependent on the assumptions). For the leptophilic ALP they are relevant, since they constrain a portion of the parameter space where ALPs are in the so-called “trapped regime,” where ALPs are too strongly coupled to the plasma to escape the star. These constraints are labeled as “low energy SN (g_{ee})” in Figs. 8 and 9 (left), as opposed to the previously mentioned cooling bounds that are labeled as “SN cooling (g_{ee}).” A similar argument was used to set limits on $g_{e\mu}$ in [107], which we also show in Figs. 5 and 6.

Beam dumps can pose meaningful constraints on the LFC case despite the smaller ALP coupling to electrons, as shown in Figs. 8 and 9 (left). In particular, the E137 [72], E141 [73], and Orsay [74] electron beam dumps have been long known to place limits on axionlike particles [70,71]. For E137, we plot the results of Ref. [81] for the flavor universal case, and Refs. [68] and [69] for the flavor-conserving muonophobic ($g_{\mu\mu} = 0$) case. We find that these improve on the limits from CHARM and SINDRUM shown in Ref. [15].

Colliders, such as *BABAR* [75] and *Belle* [76] have set limits on leptophilic scalars that can be adapted to the ALP scenarios considered here. The searches are based on $e^+e^- \rightarrow \tau^+\tau^-a(a \rightarrow \ell^+\ell^-)$ and set a lower limit on f_a of $\mathcal{O}(500 \text{ GeV})$ for $m_a > 40 \text{ MeV}$. In most of the excluded parameter space, ALPs are promptly decaying particles, and therefore, cannot be constrained by traditional beam

dumps setups. In the near future, one can also envisage new constraints on ALPs from the decays of $W^\pm$ bosons at the LHC, produced, for example, in $W \rightarrow \tau(a \rightarrow \ell^+\ell^-)\nu_\tau$ decays [108]. These can help constrain short-lived ALPs heavier than the tau up to $m_W - m_\tau$. In the tauphilic scenario, most other probes of the ALP couplings rely on its loop-induced coupling to photons. Searches for $e^+e^- \rightarrow \gamma a_{\text{inv}}$ at BES-III [109] and $e^+e^- \rightarrow \gamma(a \rightarrow \gamma\gamma)$ at BES-III [110] and Belle-II [111] constrain decay constants of several GeV, as derived in [86].

The Belle-II experiment, in particular, will have a rich program to explore new physics in tau decays. Some examples include the mono-photon signatures mentioned above, a $(g - 2)_\tau$ measurement, and searches for rare tau decays (see, for example, Ref. [112] for a review). Therefore, we can expect improvements to several of the ALP constraints discussed here in the near future.

Fixed target experiments such as NA64 can constrain ALPs appearing either as missing energy [82,113], through their decays to e^+e^- [77,83], or, in the future, through LFV processes in the target, such as $\mu \rightarrow e$ [114] or $\mu \rightarrow \tau$ [115]. Searches have been done using both an electron and a muon beam [84], shown as NA64e or NA64μ, respectively. For the flavor-conserving scenarios, we draw the contours from Ref. [81]. Because of the higher energy at fixed target experiments and the access to relatively short-lived particles, they can reach ALP masses larger than the tau mass [115–118].

Anomalous magnetic moments of the electron and muon are measured very precisely and are also sensitive to novel contributions from an ALP running in the loop. In the case of the muon, the difference between the experimental world average for $a_\mu = (g_\mu - 2)/2$, including the recent measurement by the Muon g-2 experiment [119], and the latest SM theoretical prediction [120] is given by

$$\Delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{TH}} = (384 \pm 637) \times 10^{-12}. \quad (4.9)$$

Note that the SM theory prediction is based on lattice calculations of the hadronic vacuum polarization contribution [121–124], while the dispersion based estimates are in tension with it and with the experimental value for a_μ (see the discussion in [120]). Since ALPs contribute to a_μ with a negative sign to Δa_μ , we can place a limit on the ALP contribution by requiring $\Delta a_\mu > -890 \times 10^{-12}$ at 2σ . For the flavor-universal LFC ALP benchmark, this limit excludes decay constants as large as $f_a \gtrsim \mathcal{O}(10^2) \text{ GeV}$, closing the window between meson decay and beam-dump constraints. In the case of the electron Δa_e , the experimental landscape is less conclusive, with the two leading measurements disagreeing at more than 5σ [125,126]. For that reason, in Fig. 8, we only show the limits from a_μ and not a_e . In deriving the excluded regions we use the expressions in Ref. [101], neglecting diagrams proportional

to the ALP coupling to two photons that are subdominant in our region of interest.

V. CONCLUSIONS

We investigated the phenomenology of leptophilic ALPs, with a particular focus on ALPs that couple only to leptons. These can be produced in tau decays, providing a new opportunity to discover or exclude ALPs with couplings that would otherwise be beyond the reach of flavor, beam dump, and collider observables (see, e.g., Refs. [9,14,101,127–131]). We considered both flavor-violating and flavor conserving ALP couplings to leptons, with or without flavor universality, allowing for ALP decays to be prompt, displaced/long-lived, or escape detection entirely—appearing as missing energy and momentum. We derived new bounds and estimated the sensitivity of current and future experiments to the decay constant f_a of the ALP. Ultimately, the SHiP experiment will be well poised to probe regions of parameter space with decay constants as high as $\mathcal{O}(10^9 \text{ GeV})$ for flavor-violating ALP scenarios in the mass range $m_\mu + m_e < m_a < m_\tau - m_e$.

Past searches for long-lived particles at the CHARM and BEBC beam dumps at CERN already constrain the parameter space of leptophilic ALPs, reaching decay constants as high as $f_a \sim \mathcal{O}(10^8 \text{ GeV})$ for the lepton-flavor-violating case. We have also presented the event-rate sensitivity (corresponding to 90% CL sensitivity in the case of negligible background counts) of currently operating experiments: MicroBooNE, ICARUS, and NOvA at Fermilab, and protoDUNE and FASER at CERN, as well as of future experiments such as the next-generation forward physics facility at the LHC, FASER-2, the dedicated large-acceptance beam-dump at CERN’s SPS, SHiP, and the DUNE near detector. Eventually, with its full dataset, SHiP is expected to surpass all other projects thanks to the large number of taus produced at the target and the large acceptance of the detector. The improvement over FASER and FASER-2 in the short-lived regime, where a large number of ALPs are produced but only a few travel far enough to decay inside the detector, is small. However, in the long-lived regime, where a small number of ALPs are produced but only a few decay as they pass through the detector, the improvement over other projects can be significant, with over an order of magnitude improvement over CHARM and BEBC constraints.

For flavor-violating scenarios with a hierarchy between flavor-violating and flavor-conserving couplings, $g_{\ell_i \ell_j} \ll g_{\ell_i \ell_i}$, the CHARM and BEBC exclusions are closer to $f_a \sim \mathcal{O}(10^6\text{--}10^7 \text{ GeV})$, but they do constrain new parameter space. If flavor-violating couplings are heavily suppressed, keeping the ALP production rates from $\tau \rightarrow \ell a$ decays sizable requires smaller values of f_a . In this limit, the ALP becomes very short-lived above the dimuon threshold, quickly degrading the sensitivity of conventional

beam dumps. In this scenario, FASER-2 and SHiP have an advantage since FASER-2 will be capable of producing ALP with larger boosts and SHiP will be located closer to the target.

In the case of lepton-flavor-conserving couplings, ALP production can only take place via three-body and four-body G_F -suppressed decays, and thus the experimental reach is much weaker. For a flavor universal scenario, $g_{ee} = g_{\mu\mu} = g_{\tau\tau}$, we find that FASER-2 and SHiP will be able to probe new parameter space for $f_a \sim \mathcal{O}(10 - 10^4) \text{ GeV}$. For such low values of f_a , ALPs are only long-lived, if they are lighter than $m_a < 2m_\mu$, restricting the search strategies to just the decays of ALPs to e^+e^- pairs. This conclusion can change, if the flavor-conserving couplings are nonuniversal. In particular, if $g_{\mu\mu} \ll g_{ee}, g_{\tau\tau}$, then the partial decay width for $a \rightarrow \mu^+\mu^-$ decays can be as suppressed as the partial decay width for $a \rightarrow e^+e^-$ (or even further suppressed), extending the experimental reach of beam dump experiments to higher ALP masses, above the dimuon threshold.

We conclude by emphasizing that there still remain many opportunities to probe physics beyond the Standard Model with tau leptons. We have presented one example based on long-lived particle production, the leptophilic ALP, but one may equally consider the production of other exotic short-lived states. In that case, the new-physics searches are restricted to less-intense, but cleaner environments such as e^+e^- colliders. Given the large numbers of τ leptons set to be produced at modern colliders, one can expect significant progress from Belle-II [112], LHCb, ATLAS, and CMS, and eventually at a Super Tau-Charm factory [132]. We plan to present a more detailed study of exotic signals due to short-lived dark particles produced in tau decays in a future work [133].

ACKNOWLEDGMENTS

We thank Brian Batell, Simon Knapen, Stephen Mrenna, and Maksym Ovchinnikov for helpful discussions. The work of M.H. is supported by the Neutrino Theory Network Program Grant No. DE-AC02-07CH11359 and the US DOE Award No. DE-SC0020250. A.R. acknowledges support from the National Science Foundation (Grant No. PHY-2020275) and to the Heising-Simons Foundation (Grant No. 2017-228). J.Z. and T.M. acknowledge support in part by the DOE Grant No. DE-SC0011784, and NSF Grants No. OAC-2103889, No. OAC-2411215, and No. OAC-2417682. Fermilab is managed by FermiForward Discovery Group, LLC, acting under Contract No. 89243024CSC000002 with the U.S. Department of Energy, Office of Science, Office of High Energy Physics. This research was supported by the Munich Institute for Astro-Particle and BioPhysics (MIAPP) which is funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) under Germany’s Excellence

Strategy—EXC-2094–390783311. Part of this research was performed at the Aspen Center for Physics, which is supported by National Science Foundation Grant No. PHY-1607611. This work was partially supported by the University of Iowa’s Year 2 P3 Strategic Initiatives Program through funding received for the project entitled “High Impact Hiring Initiative (HIHI): A Program to Strategically Recruit and Retain Talented Faculty.

DATA AVAILABILITY

The data that support the findings of this article are openly available [134], embargo periods may apply.

APPENDIX A: IMPACT OF TAU PRODUCTION UNCERTAINTIES

The total number of charmed mesons and their energy spectrum at fixed target experiments is subject to sizeable uncertainties. For the τ yield, the uncertainty on $D_s^\pm$

production is the most important one ($D^\pm$ have a smaller branching ratio to taus). In this appendix, we discuss how these uncertainties can impact the experimental sensitivities to leptophilic ALPs from tau decays, using NOvA and SHiP for illustration. We compare the methodology of Sec. III B, based entirely on our PYTHIA simulation, with an approach based on a simplified parametrization of the $D_s^\pm$ production rate and kinematics, following Ref. [32].

A standard parametrization for a $D^\pm$ and D^0 production cross section at fixed target experiments is [135]

$$\frac{d\sigma}{dp_T dx_F} \propto (1 - |x_F|)^n e^{-ap_T - bp_T^2}, \quad (\text{A1})$$

where p_T is the transverse momentum of the outgoing D meson and $x_F = p_z^{\text{COM}}/p_{z-\text{max}}^{\text{COM}}$ is the Feynman x calculated in the center of mass of the pp collision. In this procedure, the normalization is fixed by the charm production cross section. The number of τ leptons N_τ produced

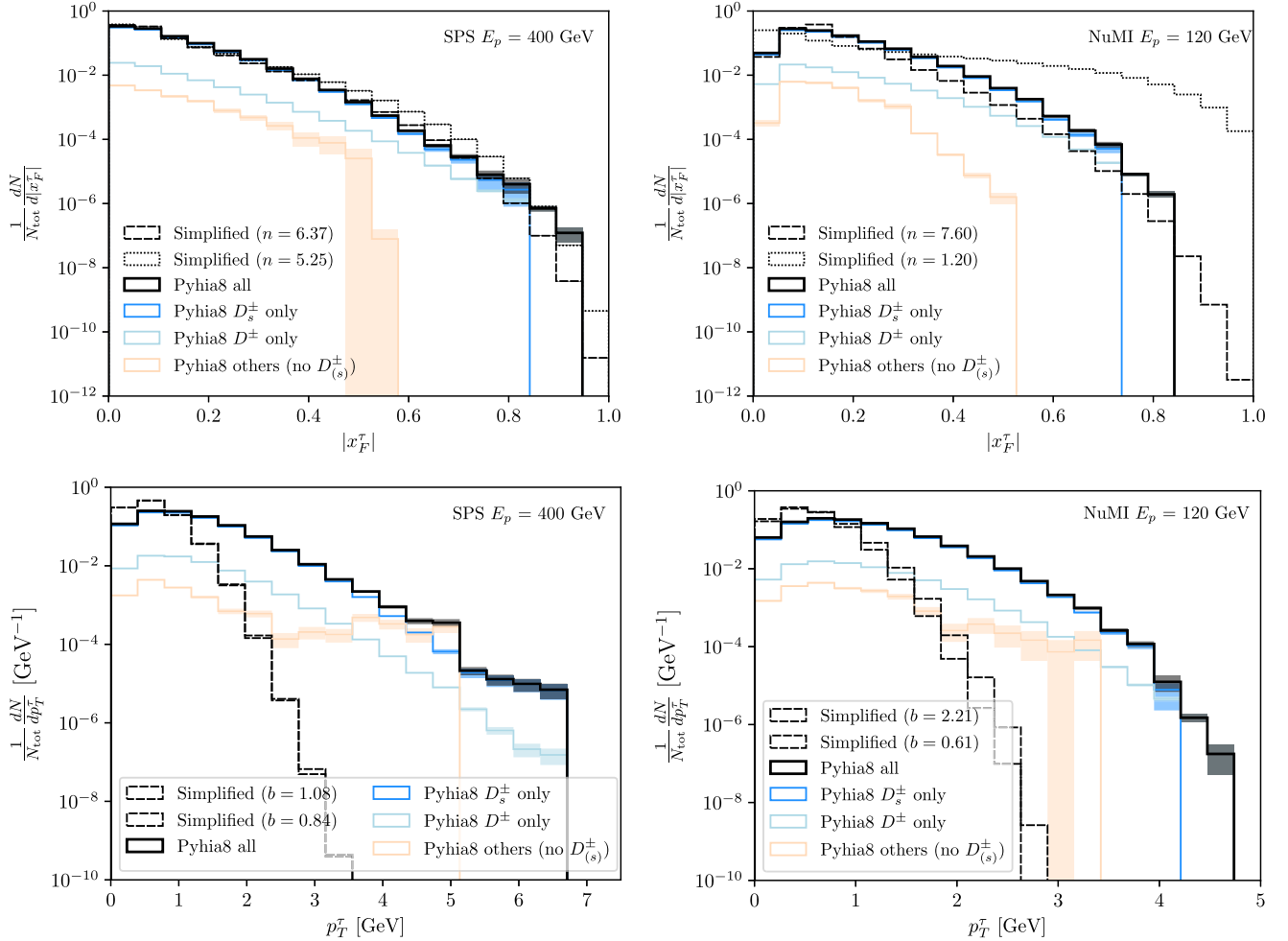


FIG. 10. x_F (top row) and p_T (bottom row) distributions of $\tau^\pm$ leptons produced in primary pp collisions at the NuMI and SPS target in our PYTHIA simulation (solid black lines) compared with the simplified treatment following the parametrization in [32] (dashed grey lines). In this simplified treatment, all taus are produced by $D_s^\pm$, which are assumed to follow the same kinematics as the $D^\pm$ of [32]. We also show our PYTHIA results including taus only from $D_s^\pm$ mesons (blue solid), $D^\pm$ mesons (light blue solid), and all other sources of taus (cream solid).

in pp collisions is approximately given by

$$N_\tau = N_{\text{pp-col}} \frac{\sigma_{c\bar{c}}}{\sigma_{pp}^{\text{inelastic}}} f_{c \rightarrow D_{(s)}} \mathcal{B}(D_s \rightarrow \tau), \quad (\text{A2})$$

where $N_{\text{pp-col}}$ is the number of proton-proton collisions, $\sigma_{c\bar{c}}$ the total charm-pair production cross section, and $\sigma_{pp}^{\text{inelastic}}$ the total inelastic pp cross section. Note that in the above expression we neglected the $D^\pm$ contribution. The fragmentation fraction is given by $f_{c \rightarrow D_s} = 0.079$ [136] and the branching ratio of all the modes containing τ set to $\mathcal{B}(D_s \rightarrow \tau) = 5.4\%$ as in the main text.

In Ref. [32], the parameters n , a , b and as well as the charm production cross section $\sigma_{c\bar{c}}$ were tuned simultaneously to a broad set of results generated from a custom version of PYTHIA that incorporates FONLL matrix elements [137,138] and used LHAPDF6 [139] and the NNPDF4.0 set of parton-distribution functions [140].

Adopting the prescription of Ref. [32], we vary these inputs in concert and propagate the resulting spread to the projected ALP coverage. While this procedure is not a rigorous estimate of the theoretical uncertainty, it illustrates how variations in the charm-production distributions can impact the experimental reach of searches for leptophilic ALPs produced from tau decays. In the absence of a dedicated study for $D_s^\pm$, we adopt the parameters and ranges quoted for $D^\pm$. In what follows, we quote the range of values used in [32],

$$\text{NuMI: } (n, b) = (4.40 \pm 1.60, 1.41 \pm 0.40), \quad (\text{A3})$$

$$\text{SPS: } (n, b) = (5.81 \pm 1.3, 0.96 \pm 0.06), \quad (\text{A4})$$

where we also set $a = 0$. We find that the number of taus per POT with our PYTHIA simulation is in agreement with the normalization of [32] using Eq. (A2). In our comparisons below we normalize the simplified simulations according to the results of the fit to $\sigma_{c\bar{c}}$.

In Fig. 10, we compare the resulting p_T and x_F distribution of taus from our PYTHIA simulations with those obtained from the simplified procedure outlined above. We plot the results of the simplified approach using the extremes of the 2σ allowed regions for the parameters in [32]. Where not specified, we use $n = 4.4$ and $b = 1.41$ for NuMI and $n = 5.81$ and $b = 0.96$ for SPS. While we find some agreement in the longitudinal variable x_F , we see a significant deviation in the p_T spectrum with respect to PYTHIA results. The p_T distribution directly affects the geometrical acceptance of different experiments to ALPs, and will thus be a source of uncertainty. The tails of the longitudinal component of the momentum can also extend to higher values in the simplified approach, impacting the largest boost achievable for ALPs and, therefore, the sensitivity of some experiments to the shorter-lived regime of the ALP parameter space.

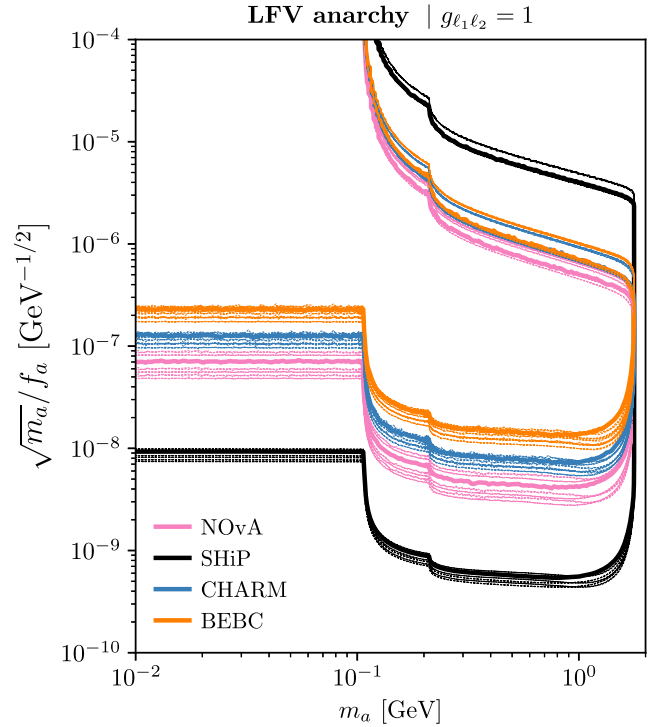


FIG. 11. Projected event-rate sensitivities of NOvA (pink), CHARM (blue), BEBC (orange), and SHiP (black) using our standard PYTHIA simulation (thick solid line) and the simplified simulation based on a parametrization of $D_s^\pm$ production with the 2σ ranges of input normalization and shape parameters from [32] (thin dashed lines). While the normalization uncertainty impacts mostly the upper limit (small f_a^{-1}), the difference between the PYTHIA results and this simplified simulation are also noticeable on the lower limit (large f_a^{-1}) due to the uncertainties on the kinematics of the $D_s^\pm$ produced and the resulting boost factor for the ALPs.

Figure 11 compares the resulting event-rate sensitivities of our PYTHIA results with the simplified approach described above (using the 2σ extreme values) in the parameter space of $(\sqrt{m_a}/f_a)$ versus m_a as some parts of the sensitivity to $\sqrt{m_a}/f_a$ are flat in m_a , as opposed to the sensitivity to $1/f_a$. We see that while the normalization impacts mostly the upper limit, the PYTHIA simulation and the simplified approach also disagree in the short-lived regime (lower limit) due to the different kinematics of the ALPs.

In the future, perturbative QCD calculations of the double-differential production rates of D_s mesons at the NuMI and SPS targets could refine the τ production calculation presented here, similarly to what was performed for for SHiP in Ref. [25,26] and for FASER in Refs. [141–144].

APPENDIX B: PARAMETRIZED TAU PRODUCTION

In the main text, we describe our simulation using PYTHIA events. In this appendix, we present an alternative

method for fast simulation of tau production at NuMI, SPS, and the LHC. We do not rely on this method in our results, but include it here for future reference. In analogy to parametrizations historically used for charmed hadron production, we can describe the kinematics of the *tau* lepton in terms of its transverse momentum p_T and $x_F = p_z^{\text{COM}}/p_{z-\text{max}}^{\text{COM}}$, the Feynman- x of the tau calculated in the center of mass of the pp collision. We find that, up to an overall normalization factor, a satisfactory parametrization of the tau kinematics at the three experiments can be found with the following exponential functions:

$$\frac{dN_\tau}{dx_F dp_T} = \left[\sum_{i=1}^3 g_i \Gamma(p_T, \alpha_i, \lambda_i) \right] \left[\sum_{j=1}^3 c_j \exp(-a_j |x_F|^{n_j}) \right], \quad (\text{B1})$$

where $\Gamma(x, \alpha, \lambda) = x^{\alpha-1} e^{-\lambda x} \lambda^\alpha / \Gamma(\alpha)$ is the gamma distribution and $\Gamma(\alpha)$ the gamma function. Here c_i , g_i , n_i , α_i , and λ_i

are all constants that we fit to the output of our PYTHIA simulations. As expected, we find that x_F and p_T are, to a good approximation, uncorrelated, so our fits can be done separately in terms of the two one-dimensional distributions.

The resulting best-fit curves are shown alongside the PYTHIA simulation output in Fig. 12 and the best-fit parameters are listed in Table II. In $x_F(p_T)$, we restrict the range of the fit to $|x_F| < 1$, and $p_T < 5$ GeV for NuMI, $p_T < 6$ GeV for SPS, and $p_T < 30$ GeV for the LHC. With x_F and p_T distributions fixed in the center of mass of the pp collision, one can sample the remaining azimuthal angular variable in a uniform way and boost the system to the appropriate lab frame.

APPENDIX C: DETAILS ON ALP PRODUCTION AND DECAY

In this appendix, we summarize the computational details of the ALP production and decay rates.

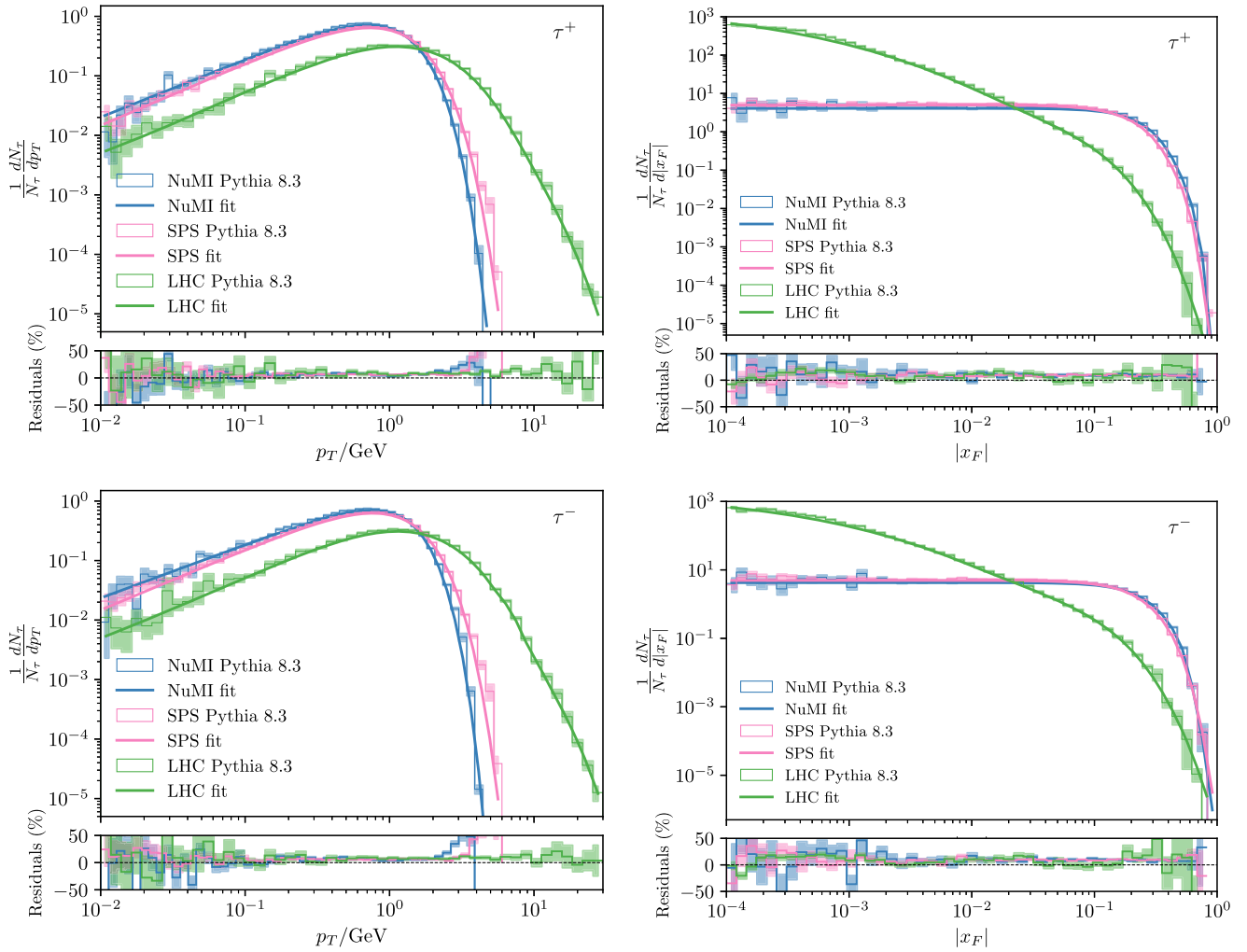


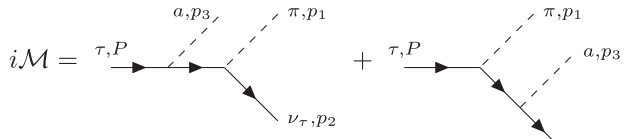
FIG. 12. $|p_T|$ (left) and $|x_F|$ (right) distributions of τ^+ mesons from primary pp collisions at the NuMI and SPS targets as well as from the LHC. The results from our PYTHIA simulations (solid histograms) are fit using analytic functions (see text for details). The best fit curves are shown as solid lines. The bottom panel shows the relative residuals between simulation and our fit.

TABLE II. Best-fit parameters for τ^+ and τ^- production in pp collisions, grouped by transverse momentum (p_T) and Feynman- x (x_F) distributions. The fits consist of a sum of three gamma distributions in p_T and two exponential terms in x_F . Sets NuMI, SPS, and LHC correspond to different collider setups.

	NuMI	SPS	LHC	NuMI	SPS	LHC
Parameter	τ^+			τ^-		
Gamma distribution fits in p_T						
g_1	2047052.33	2869884.82	29.95	84411.01	9496674.67	22.30
m_1	1.684	1.440	0.857	1.896	1.511	0.850
μ_1	1.578	1.729	2.682	1.473	1.674	2.839
λ_1	0.964	0.875	0.660	1.219	0.990	0.627
g_2	0.570	0.570	4.515	0.570	0.570	3.256
m_2	1.000	1.000	0.695	1.000	1.000	0.588
μ_2	89.18	89.18	2.105	89.18	89.18	2.555
λ_2	84.40	84.40	0.838	84.40	84.40	0.535
g_3	0.570	0.570	0.109	0.570	0.570	4.628
m_3	1.000	1.000	6.769	1.000	1.000	2.166
μ_3	89.18	89.18	116.25	89.18	89.18	99.912
λ_3	84.40	84.40	0.00347	84.40	84.40	0.662
Exponential fits in x_F						
r_1	48.69	0.339	263.92	5112.29	3810.67	351.77
a_1	16.95	31519936.65	17.36	18.32	16.70	17.10
n_1	4.559	35.98	0.808	2.026	1.795	0.777
r_2	3593.68	189.77	132431.87	82.01	137.60	126552.23
a_2	18.39	17.51	23.68	119.43	2254.07	26.02
n_2	2.090	1.807	0.321	6.873	4.767	0.342

1. The $\tau \rightarrow \pi\nu_\tau a$ decay

We begin with the decay mode $\tau^- \rightarrow \pi^- \nu_\tau a$. The amplitude of this process is given by



$$i\mathcal{M} = \tau, P \rightarrow \pi, p_1 + \nu_\tau, p_2 + a, p_3 + \tau, P \rightarrow \pi, p_1 + \nu_\tau, p_2 + a, p_3$$

$$= -\frac{\sqrt{2}iG_F V_{ud} f_\pi m_\tau g_{\tau\tau}}{f_a(m_\tau^2 - \bar{m}_\tau^2)} \bar{u}_\nu P_R [m_\tau \not{p}_{12} - \bar{m}_\tau^2] u_\tau, \quad (C1)$$

where $P_R = (1 + \gamma_5)/2$, $p_{ij} = p_i + p_j$, and $\bar{m}_\tau^2 = p_{12}^2$ is the invariant mass of the intermediate off-shell τ (in the first diagram). We need to include the diagram with the ALP attached to the neutrino leg to be consistent with the electroweak gauge symmetry; only after this, we obtain the amplitude proportional to m_ℓ .

The differential decay rate is given by

$$d\Gamma(\tau \rightarrow \pi\nu_\tau a) = \frac{|\mathcal{M}|^2}{2m_\tau} d\Phi_3, \quad (C2)$$

where $d\Phi_3$ is the three-body phase space element [with the $(2\pi)^4$ included in the delta function that ensures the total

four-momentum conservation]. Averaging over the spin of the initial tau, gives

$$d\Gamma(\tau \rightarrow \pi\nu_\tau a) = \frac{|\bar{\mathcal{M}}|^2}{256\pi^3 m_\tau^3} d\bar{m}_\tau^2 dm_{23}^2, \quad (C3)$$

where $|\bar{\mathcal{M}}|^2$ is the spin averaged amplitude squared and $m_{23}^2 = p_{23}^2$. Integrating this expression over m_{23}^2 from $(m_{23}^2)_+$ to $(m_{23}^2)_-$, where

$$(m_{23}^2)_\pm = (E_2^* + E_3^*)^2 - \left(E_2^* \pm \sqrt{E_3^{*2} - m_a^2}\right)^2, \quad (C4)$$

and

$$E_2^* = \frac{\bar{m}_\tau^2 - m_\pi^2}{2\bar{m}_\tau}, \quad E_3^* = \frac{m_\tau^2 - \bar{m}_\tau^2 - m_a^2}{2\bar{m}_\tau}, \quad (C5)$$

we obtain Eq. (2.5) in the main text.

2. The $\tau \rightarrow \rho\nu_\tau a$ decay

The calculation of partial decay width for $\tau^- \rightarrow \rho^- \nu_\tau a$, is analogous to the one for $\tau^- \rightarrow \pi^- \nu_\tau a$. The amplitude is given by

$$i\mathcal{M} = \tau, P \rightarrow \ell, p_1 + \nu_\tau, p_2 + a, p_3 + \nu_\tau, p_2$$

$$= -\frac{\sqrt{2}G_F V_{ud} m_\rho m_\tau f_\rho g_{\tau\tau}}{f_a(m_\tau^2 - \bar{m}_\tau^2)} \bar{u}_\nu P_R \not{\epsilon}^* [\not{p}_{12} - m_\tau] u_\tau, \quad (\text{C6})$$

where we again note the contribution from the second diagram. After performing the phase space integrals, averaging over the tau spin, and summing over final state particle polarizations, we obtain Eq. (2.9).

3. The $\tau \rightarrow \ell \nu_\ell \nu_\tau a$ decay

The calculation of the partial decay width for $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau a$, is more complicated, since it involves integrating over the four-body phase space. The amplitude is given by

$$i\mathcal{M} = \tau, P \rightarrow \ell, p_1 + \nu_\tau, p_3 + \bar{\nu}_\ell, p_4 + \dots$$

$$= -\frac{2\sqrt{2}G_F g_{\tau\tau} m_\tau}{f_a(p_{234}^2 - m_\tau^2)} [\bar{u}_\nu P_R \gamma_\alpha (\not{p}_{234} - m_\tau) u_\tau] [\bar{u}_\ell \gamma^\alpha P_L v_\nu], \quad (\text{C7})$$

where $P_L = (1 - \gamma_5)/2$, $p_{234} = p_2 + p_3 + p_4$, and for brevity we show only the diagram with the ALP attached to τ and not the one with ALP attaching to ν_τ line. In the calculation we do ignore the diagrams with the ALP attached to ℓ and $\bar{\nu}_\ell$ since those are suppressed by m_ℓ/m_τ at the amplitude level. The decay rate is given by

$$d\Gamma = \frac{1}{2m_\tau} d\Phi_4 |\bar{\mathcal{M}}|^2, \quad (\text{C8})$$

where $d\Phi_4$ is the differential four-body phase space element. The spin averaged amplitude squared is given by

$$|\bar{\mathcal{M}}|^2 = \frac{64G_F^2 |g_{\tau\tau}|^2 m_\tau^2}{f_a^2(m_\tau^2 - p_{234}^2)^2} (p_2 \cdot p_3) p_4^\alpha \times [(m_\tau^2 - p_{234}^2) P_\alpha + 2p_{234\alpha}(P \cdot p_{234} - m_\tau^2)]. \quad (\text{C9})$$

It depends on p_2 , p_3 , and p_4 separately only through the form $(p_2 \cdot p_3) p_4^\alpha$. It is then useful to decompose the four-body phase space integral sequentially as

$$\int d\Phi_4 = \int_{m_\tau^2}^{(m_\tau - m_a)^2} \frac{d\bar{m}_\tau^2}{2\pi} \int d\Phi_{1;234} \times \int_0^{(\bar{m}_\tau - m_l)^2} \frac{d\bar{m}_{\nu\nu}^2}{2\pi} \int d\Phi_{2;34} \int d\Phi_{3;4} \quad (\text{C10})$$

where $\bar{m}_\tau^2 = p_{234}^2$ and $\bar{m}_{\nu\nu}^2 = p_{34}^2$ are the off-shell invariant mass of τ and $\bar{\nu}_\ell \nu_\tau$ pair, respectively, and

$$\int d\Phi_{1;234} = \frac{1}{4\pi^2} \int \frac{d^3 p_1}{2E_1} \frac{d^3 p_{234}}{2\bar{E}_{234}} \delta^{(4)}(P - p_1 - p_{234}), \quad (\text{C11})$$

$$\int d\Phi_{2;34} = \frac{1}{4\pi^2} \int \frac{d^3 p_2}{2E_2} \frac{d^3 p_{34}}{2\bar{E}_{34}} \delta^{(4)}(p_{234} - p_2 - p_{34}), \quad (\text{C12})$$

$$\int d\Phi_{3;4} = \frac{1}{4\pi^2} \int \frac{d^3 p_3}{2p_3} \frac{d^3 p_4}{2p_4} \delta^{(4)}(p_{34} - p_3 - p_4), \quad (\text{C13})$$

with $\bar{E}_{234}^2 = |\vec{p}_{234}|^2 + \bar{m}_\tau^2$ and $\bar{E}_{34}^2 = |\vec{p}_{34}|^2 + \bar{m}_{\nu\nu}^2$. We can show, e.g., by going to the p_{34}/p_{234} rest frame and using the Lorentz invariance, that

$$\int d\Phi_{3;4} p_3^\alpha p_4^\beta = \frac{1}{96\pi} (\eta^{\alpha\beta} \bar{m}_{\nu\nu}^2 + 2p_{34}^\alpha p_{34}^\beta), \quad (\text{C14})$$

and therefore

$$\int d\Phi_{2;34} \int d\Phi_{3;4} (p_2 \cdot p_3) p_4^\alpha = \frac{I(\bar{m}_\tau, \bar{m}_{\nu\nu})}{1536\pi^2} \bar{m}_\tau^2 p_{234}^\alpha, \quad (\text{C15})$$

where the function $I(\bar{m}_\tau, \bar{m}_{\nu\nu})$ is defined in Eq. (2.13). By using

$$P \cdot p_{234} = \frac{1}{2} [m_\tau^2 + \bar{m}_\tau^2 - m_a^2], \quad (\text{C16})$$

and performing the angular integral of the “1; 234” system, we obtain Eq. (2.12).

4. The $a \rightarrow \gamma\gamma$ decay

Finally, we consider the ALP decay mode $a \rightarrow \gamma\gamma$ through the loop of a heavy lepton ℓ . This process is identical to the chiral anomaly in the limit $m_\ell \rightarrow 0$, and therefore we need some care on the regularization. The amplitude is diagrammatically expressed as

$$i\mathcal{M} = a, P \rightarrow \gamma, p_1 + \gamma, p_2 \quad (\text{C17})$$

Its evaluation is similar to the chiral anomaly computation, and we obtain

$$i\mathcal{M} = \frac{2i\alpha}{\pi} \epsilon^{\alpha\beta\mu\nu} p_{1\mu} p_{2\nu} \epsilon_\alpha^*(p_1) \epsilon_\beta^*(p_2) \times \int_0^1 dx \int_0^{1-x} dy \left[1 - \frac{m_\ell^2}{m_\ell^2 - xym_a^2} \right]. \quad (\text{C18})$$

This expression reproduces the chiral anomaly in the limit $m_\ell \ll m_a$, while it vanishes in the opposite limit $m_\ell \gg m_a$. In general, the anomaly equation for a massive fermion contains two contributions; $\partial_\mu(\bar{\ell}\gamma^\mu\gamma_5\ell)$ and $m_\ell\bar{\ell}i\gamma_5\ell$. In the high energy (or massless) limit, the anomaly equation,

i.e., the $F\tilde{F}$ operator, is saturated by the former contribution, while in the low energy (or heavy mass) limit, it is saturated by the latter contribution. One can perform the Feynman parameter integrals analytically, leading to Eq. (2.18) after the phase space integration.

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