TABLE II. Explicit comparison with experimental $^3P^e$–$^3D^0$ transitions.

	λ (Å)	
Level(n)	Calculated	Observed (vac.)
1	3014.46	3014.6±3
2	2562.71	2561.76±1
3	2400.91	2402.71±1
4	2321.91	2319.69±1

pansion results of Doyle *et al.*[3] Careful optimization of the energy with respect to four nonlinear parameters still did not allow us to approach the $1/z$ results and we believe those results are too low.

The mass-polarization correction

$$E_M = -(2/M)\langle \Phi \vec{\nabla}_1 \cdot \vec{\nabla}_2 \Phi \rangle$$

is given for all the states in Table I. Finally, the positions of the D states, including the reduced mass and the mass-polarization corrections, with respect to the ground state of He are given in Table I.

Combining the $^3D^0(1)$ results with the position of $^3P^e$ state,[5] we find the wavelength of the transition $^3P^e \rightarrow {}^3D^0$ is equal to 3014.16 Å. This agrees with the results of Doyle *et al.*[3] and also with the recent experimental results of Berry *et al.*[6] An explicit comparison with the experimental $^3p^e - {}^3D^0$ transitions is given in Table II.

No $^{1,3}D^0$ states in H^- were found in this calculation.

I wish to thank Dr. A. Temkin for critical reading of the manuscript.

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ERRATUM

Electronic Transmission Spectroscopy: Core-Excited Resonances in Diatomic Molecules, L. Sanche and G. J. Schulz [Phys. Rev. A 6, 69 (1972)]. Figure 10 has the band designations incorrectly labeled. Below we show the corrected Fig. 10.

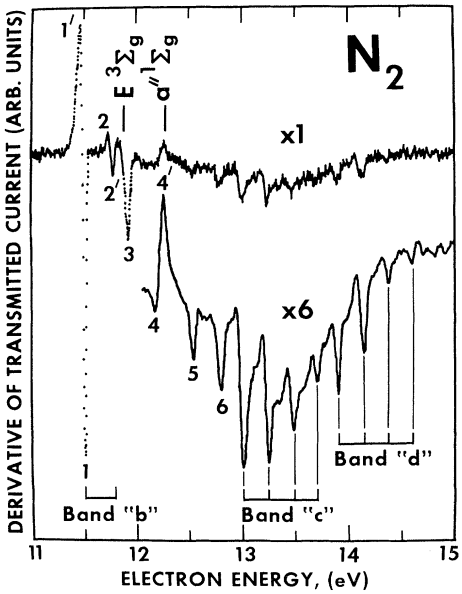

FIG. 10. Derivative of transmitted current vs electron energy in N_2. The giant resonance $1-1'$ is a $^2\Sigma_g^+$, N_2^- state whose parents are the $E\,^3\Sigma_g^+$ and $a''\,\Sigma_g^+$ Rydberg states of N_2. The grandparent of the structures 1–4 is the ground state of N_2^+. The other resonances, including bands "c" and "d" which appear on the higher-sensitivity run at the bottom of the figure, have the $A\,^2\Pi_u$ state of N_2^+ as grandparent.